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TESTING THE PARADIGM THAT ULTRALUMINOUS X-RAY SOURCES AS A CLASS REPRESENT ACCRETING INTERMEDIATE-MASS BLACK HOLES

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ABSTRACT

To test the idea that ultraluminous X-ray sources (ULXs) in external galaxies represent a class of accreting intermediate-mass black holes (IMBHs), we have undertaken a program to identify ULXs and a lower luminosity X-ray comparison sample with the highest quality data in the *Chandra* archive. We establish as a general property of ULXs that the most X-ray-luminous objects possess the flattest X-ray spectra (in the *Chandra* bandpass). No prior sample studies have established the general hardening of ULX spectra with luminosity. This hardening occurs at the highest luminosities (absorbed luminosity $\geq 5 \times 10^{39}$ erg s $^{-1}$) and is in line with recent models arguing that ULXs are actually stellar mass black holes. From spectral modeling, we show that the evidence originally taken to mean that ULXs are IMBHs—i.e., the “simple IMBH model”—is nowhere near as compelling when a large sample of ULXs is looked at properly. During the last couple of years, *XMM-Newton* spectroscopy of ULXs has to a large extent begun to negate the simple IMBH model based on fewer objects. We confirm and expand these results, which validates the *XMM-Newton* work in a broader sense with independent X-ray data. We find that (1) cool-disk components are present with roughly equal probability and total flux fraction for any given ULX, regardless of luminosity, and (2) cool-disk components extend below the standard ULX luminosity cutoff of 10^{39} erg s $^{-1}$, down to our sample limit of $10^{38.3}$ erg s $^{-1}$. The fact that cool-disk components are not correlated with luminosity damages the argument that cool disks indicate IMBHs in ULXs, for which strong statistical support was never found.

Subject headings: accretion, accretion disks — galaxies: general — surveys — X-rays: binaries

Online material: color figures

1. INTRODUCTION

Ultraluminous X-ray sources (ULXs) have long been hailed as direct observational evidence for the existence of accreting intermediate-mass black holes (IMBHs; Colbert & Mushotzky 1999). The X-ray spectral model that has emerged as a central pillar for this argument (the “simple IMBH model”) is that which is commonly applied as the canonical X-ray spectral fit to Galactic black hole binaries with stellar mass black holes (McClintock & Remillard 2006). This model consists of a thermal accretion disk component plus a power-law (PL) continuum component. When applied to ULX spectra, the derived disk temperatures are 0.1–0.3 keV (e.g., Miller et al. 2004), much lower than Galactic black holes (at 0.6–1 keV). A cooler disk implies a bigger disk; so, assuming that the disk approximately extends inward to the last stable orbit around the black hole, this would imply bigger and more massive black holes. Such cool disks were indeed found in a few ULXs (e.g., Miller et al. 2004).

To counter this argument, many recent papers have pointed out both theoretical and observational problems with the simple IMBH model as a global explanation for all ULXs (e.g., Gonçalves & Soria 2006; Stobbart et al. 2006b; Roberts 2007). The observed accretion disk components can be fairly weak; thus, they do not provide a reliable measure of black hole mass. Also, the simple

IMBH model does not necessarily approximate well the X-ray spectra of many ULXs. Attention has switched to perhaps less exotic models to explain some of the ULXs, such as beaming (King et al. 2001) or super-Eddington accretion (Begelman 2002), both of which explain ULX X-ray properties without the need for an IMBH. Galactic super-Eddington sources are known, such as stellar mass black hole binaries like GRS 1915+105 (Fender & Belloni 2004), V4641 Sgr (Revnivtsev et al. 2002), and possibly SS 433. The latter could be an example of both beaming *and* super-Eddington emission, the combination of which could easily explain even the most luminous ULXs (Begelman et al. 2006; Poutanen et al. 2007). Cool accretion disks can also be physically explained by “coupled disk-corona” models (Done & Kubota 2006), blurred emission and absorption lines from surrounding (outflowing) gas (Gonçalves & Soria 2006), or a microblazar with magnetized jets (Freeland et al. 2006) that can transfer disk energy into the jet (thus making the disk fainter and cooling it at the same time).

Recent detailed X-ray spectral modeling has revealed properties that further complicate any simple global interpretation, suggesting that multiple classes of ULXs exist. Some very bright ULXs have been found by several authors (Zezas et al. 2002; Soria et al. 2007; Soria & Wong 2006; Socrates & Davis 2006) to have relatively flat spectra, not usually expected in high states for

accretion states of black holes (McClintock & Remillard 2006). A flat spectrum suggests an inverse correlation between the slope of the spectrum and source luminosity (see also NGC 5204 X-1; Roberts et al. 2006). Such an inverse correlation is hard to explain with current IMBH models, because in the typical high state the spectrum is soft, dominated by the disk component and with a steep PL (McClintock & Remillard 2006). Specifically, *XMM-Newton* spectroscopy of ULXs has to a large extent already begun to directly negate the simple IMBH model (see, e.g., Stobbart et al. 2006b; Gonçalves & Soria 2006).

Two *Chandra* surveys suggest that ULXs may in fact be an extension of normal lower luminosity galaxy populations to higher luminosities. Using simple PL models applied to spectra with typically 50 counts each, Swartz et al. (2004) compared ULXs to a lower luminosity sample of X-ray sources with $L_X = 10^{38} - 10^{39} \text{ erg s}^{-1}$ and found both distributions of photon indices to be well fitted by Gaussians centered at about 1.9. The samples also have similar X-ray colors, time series, and positions within their host galaxies. In another *Chandra* study, Colbert et al. (2004) found no discernible difference between the X-ray colors of ULXs and lower luminosity sources in spiral galaxies. Both analyses were done with data of fairly poor spectral quality in terms of fitting detailed models, and the latter relied on a color-color analysis rather than spectral fitting. These works also did not include two-component model fits that would identify spectral states and directly test the simple IMBH model.

To search carefully for spectral properties that can differentiate ULXs, we need the best data available that will allow us to distinguish between simple spectral models. In this paper we use the highest quality X-ray spectra for a large, complete sample of ULXs from the *Chandra* archives to test various ULX models. We are able to provide a statistically strong comparison of the results with lower luminosity X-ray sources of equal data quality. We pay special attention to properties usually associated with ULXs, such as the signature of a cool disk, which has never been tested for a uniform and large sample of good-quality *Chandra* spectra. Is the cool disk preferentially found in ULXs? If yes, does the disk dominate the total emission? We also search for other spectral behaviors found more recently in individual ULXs, such as a correlation between hardness and luminosity, and what this might mean. For the first time, strong statistical tests of various ideas of ULX models can be provided to the ULX community.

With its unmatched spatial resolution, *Chandra* is better suited than *XMM-Newton* for studying point sources in crowded regions or resolving point sources in distant galaxies. This is particularly true for the starburst galaxies that host populations of ULXs (e.g., NGC 3256, Lira et al. 2002; Cartwheel galaxy, Gao et al. 2003), where only *Chandra*'s unparalleled X-ray optics can spatially and spectrally resolve the emission of ULXs from that of the underlying galaxy. We have searched all public data available in the *Chandra* archive for ULXs and lower luminosity comparison objects with at least 1000 counts. In § 2 we present our source selection process, methods for identifying rejected objects, and an estimate of contamination from background objects. In § 3 we discuss the spectral fitting procedures and compare the spectral properties of the two samples. Our goal is to determine whether ULXs as a class have different spectral properties than the less luminous, “normal” X-ray sources and to offer an improved diagnosis by using the high-quality spectral data available in the *Chandra* archive. In § 4 we present results from the variability analysis. Finally, in § 5 we interpret our results and discuss the insight provided into the nature of ULXs.

2. SAMPLE SELECTION AND OBSERVATIONS

2.1. Sample Criteria

There are many published papers that address the nature of ULXs. These analyses are typically drawn from heterogeneously selected samples, small numbers of objects, or large samples with limited data quality (Swartz et al. 2004; Colbert et al. 2004). Comparisons of ULXs with other types of X-ray point sources in nearby galaxies often use selection criteria that do not provide the spectral data quality that allows a robust set of statistical conclusions to be drawn from data modeling. In this work, we use criteria that create the best possible available sample to address the nature of ULXs by defining a large and statistically robust sample of ULXs and other pointlike X-ray sources in nearby galaxies with uniform data quality. Uniformity of data quality is our prime objective, and the completeness of our sample is limited by the observations that are available in the *Chandra* archive, most of which have been obtained by other researchers for a variety of purposes.

Our ULX and comparison samples are statistically robust in the sense that we include all point sources in the *Chandra* archive with at least 1000 counts and a luminosity above $10^{38.3} \text{ erg s}^{-1}$. We carefully reject sources associated with active galactic nuclei, supernovae, and foreground stars. We also reject piled-up observations to simplify our spectral analysis. *Chandra* provides the most accurate X-ray positions to date; thus, we can be sure to identify well-isolated objects for our study. Several *XMM-Newton* studies of ULXs are published, but while these individual spectra are of higher quality, there are fewer individual point sources available due to *XMM-Newton*'s poorer imaging resolution and source confusion for faint targets located in crowded regions in galaxies. We have not made our sample fully representative in the sense of picking the same number of ULXs and comparison objects from similar galaxy types. Conclusions about the distribution of objects according to galaxy type can, however, be inferred from our statistical comparisons. Uniformly good-quality X-ray spectra allow us to apply exactly the same physical models to the ULX and comparison samples and directly compare results within the sensitivity limits.

There are selection biases inherent in our analysis. One is distance. For sources that are intrinsically less luminous, a larger fraction of objects will be located in the nearest host galaxies, while more luminous objects can be utilized from galaxies at greater distances. We also do not select objects according to any specific requirement of their local environments (e.g., their locations in their host galaxies).

2.2. Initial Sample

Our sample is derived from the list of X-ray point sources generated by the XASSIST¹ *Chandra* pipeline. For manageability, we have chosen all *Chandra* ACIS sources in the public archives as of a cutoff date of 2004 October 18. We determine which XASSIST sources are associated with host galaxies following the procedure used by Colbert & Ptak (2002). X-ray sources are further considered if they are located inside the D_{25} ellipse of their host galaxy. Parameters for the D_{25} ellipse are obtained from version 3.9b of the Third Reference Catalog of Bright Galaxies (RC3; de Vaucouleurs et al. 1991). We also only consider RC3 galaxies with recessional velocities $cz \leq 5000 \text{ km s}^{-1}$.

¹ XASSIST (A. Ptak & R. Griffiths 2003) is a semiautomatic X-ray analysis program written and maintained by A. Ptak. Analysis of archival data processed by XASSIST can be found at <http://www.xassist.org>.

To estimate observed X-ray luminosities L_{XA} , we calculated the 0.3–8.0 keV fluxes with XASSIST assuming a PL model with $\Gamma = 1.8$ and Galactic absorption and used distances for the associated RC3 galaxies. For galaxies with $cz < 1000 \text{ km s}^{-1}$, distances were taken from Tully (1988), otherwise distances were computed using $H_0 = 75 \text{ km s}^{-1} \text{ Mpc}^{-1}$. We retained all sources with $L_{\text{XA}} > 10^{38.3} \text{ erg s}^{-1}$ and then manually inspected the X-ray images to eliminate false X-ray sources chosen by the automatic data processing. This initial selection yielded 126 unique X-ray point sources in 188 *Chandra* observations. A significant fraction of the point sources were observed multiple times, which provides some useful variability information.

2.3. Obvious Rejected Objects (AGNs, QSOs, SNe, Stars, and Jets)

To reject X-ray point sources unrelated to our science, we used the NASA/IPAC Extragalactic Database (NED).² The absolute positional uncertainty for *Chandra* ACIS images is better than $1''$ (e.g., Weisskopf et al. 2003), which provides the accuracy required to identify an optical, infrared, or radio counterpart. Optical positions provided by NED are typically accurate to within a few arcseconds, and positions may be slightly less accurate for infrared and radio sources, so we first searched NED using a radius of $5''$ surrounding the XASSIST position. We next visually inspected the X-ray sources and their possible NED counterparts by overlaying the XASSIST position, the NED position, and the D_{25} galaxy ellipse onto the raw X-ray images and DSS2³ red images.

Optical images were used to check for bright foreground star counterparts. We then refined our identification search by examining the literature for more accurate positions for identified NED counterparts. In some cases VLBI measurements are available with submilliarcsecond positional accuracy, such as those used by Ma et al. (1998) for the International Celestial Reference Frame (ICRF). Some published 2MASS positions also use the ICRF reference system and have accuracies better than $0.1''$, varying slightly with the source brightness (see UCAC2; Zacharias et al. 2004). Optical positions that can be correlated with radio measurements show systematic differences of only $0.1''$ (e.g., Argyle & Eldridge 1990). Overall, we determined that the positional uncertainties of identified counterparts are generally much smaller than our X-ray positional uncertainties, the largest uncertainty being $1''$. We estimated a conservative upper limit of $1.5''$ for the net uncertainty in separation between the *Chandra* X-ray source and an identified optical, IR, or radio counterpart for any object in our sample. Therefore, we feel confident that we have identified correct counterparts to within the errors provided by the X-ray data.

From this search we reject 32 X-ray sources out of 69 *Chandra* observations. Most are associated with Seyfert and LINER galaxies (Ho et al. 1997; Véron-Cetty & Véron 2003; Bryant & Hunstead 1999). Others include background quasars and pointlike X-ray knots associated with jets within the host galaxy. As an example, source 37 in Zezas et al. (2002), in the Antennae galaxy pair (NGC 4038/4039), is a background quasar with redshift 0.26 (Clark et al. 2005). We identified a supernova in NGC 891 (SN 1986J; Bietenholz et al. 2002). One ULX in M101 (NGC 5457 X-6; Roberts & Warwick 2000) is actually a foreground star,

GSC 2.2 3842. After rejecting sources based on optical and NED counterparts, we are left with 94 X-ray sources in 119 *Chandra* observations.

2.4. Reprocessing of Archival Data and Final Rejection Criteria

Having narrowed our sample according to the above criteria, the ACIS imaging data were retrieved from the *Chandra* archives. The level 1 event files were reprocessed with CIAO version 3.0.1 and CALDB version 1.4 using the `acis_process_events` tool. No adjustment was made for charge transfer inefficiency (CTI) effects between pixels during the data readout. This allows the analysis of data to be uniform for different CCD detectors and degrees of CCD pileup. To minimize pileup effects, we restricted the count rates for on-axis full-frame (frame time 3.24 s) CCD observations to be $< 0.08 \text{ s}^{-1}$. According to the *Chandra* Proposers' Observatory Guide,⁴ this corresponds to 10% pileup. Count rates in excess of this value for point sources are likely to impact the extracted spectra. Pulse-height analyzer randomization was applied, but pixel randomization was not.

In cases where the X-ray sources were observed off-axis or in a subarray CCD mode, the pileup effect is reduced, and we can accept a higher net count rate. The actual pileup fraction is estimated for the “reduced” count rates in Table 1. These count rates were calculated by taking into account the larger point-spread function for sources observed off-axis and the CCD observation mode. Column (8) of the table lists the ACIS CCD in which the source is imaged and the subarray value, i.e., the fraction of the CCD used in the observation. Exposing a smaller chip area results in shorter frame times and reduces the pileup.

Source spectra were typically extracted from regions of radius $2''$, and local background spectra were extracted using annuli with inner and outer radii of $6''$ and $10''$. For off-axis sources we used elliptical regions, and for crowded regions, slightly more complicated background regions, as needed. Visual inspection ensured that there was no confusion with any nearby X-ray sources. Sources were retained that had > 1000 counts in the reprocessed data. Spectral fitting was performed using XSPEC version 11.2.0bd.

A total of 21 sources had less than 1000 counts after the archival data were fully reprocessed, so these are rejected. In addition, nine observations of seven sources have $> 10\%$ pileup and are rejected (Table 1). Most of the sources with significant pileup have other *Chandra* observations, so only three unique objects are fully rejected from our sample because of pileup: two ULXs and one lower luminosity source.

2.5. Final Sample

A total of 69 unique objects in 89 data sets comprise our final sample. The properties of these objects are listed in Table 2, together with some properties of their host galaxies. Using count rates derived from our reprocessed data, we recomputed the 0.3–8.0 keV observed luminosities (L_X) using a PL model with $\Gamma = 1.8$ and Galactic absorption. A final division into two groups is made according to the maximum observed luminosity, L_X^{max} . There are 47 ULXs ($L_X^{\text{max}} \geq 10^{39.0} \text{ erg s}^{-1}$) and 22 comparison objects of lower luminosity ($L_X^{\text{max}} < 10^{39.0} \text{ erg s}^{-1}$).

Some sources show luminosity variability. For two ULXs, U2 (M33 X-8) and U41 (IXO 83), their luminosity can fall below our threshold value of $10^{39.0} \text{ erg s}^{-1}$ in some cases, but we still retain the classification of ULX. Our method identifies a ULX as such if it is observed with $L_X \geq 10^{39.0} \text{ erg s}^{-1}$ at least once. On

² Available at <http://nedwww.ipac.caltech.edu>.

³ The Second Digitized Sky Survey consists of high-resolution scans of several plate collections in the red, blue, visible, and near-infrared. The images were downloaded from the server installed at ESO, using a remote client, the ESO/ST-ECF Digitized Sky Survey application.

⁴ Available at <http://cxc.harvard.edu/proposer/POG/>.

TABLE 1
PILED-UP OBSERVATIONS

Obs. (1)	Name (2)	Position (3)	Galaxy (4)	$\log L_X$ (5)	ObsID (6)	Date (7)	CCD (8)	θ_{off} (9)	Count Rate (10)	Reduced Rate (11)	Pileup (12)	Alternate Names (13)	Refs. (14)
ULXs													
1.....	U5	X031820.0–662911	NGC 1313	39.47	2950	2002 Oct 13	7(1/1)	2.4	0.25	0.13	0.15		
2.....	...	X081929.0+704219	Holmberg II	39.91	1564	2001 Nov 2	7(1/4)	0.6	0.51	0.18	0.21	IXO 31, ULX1, X-1	1, 2, 3
3.....	U10	X095550.0+694046	M82	39.52	1302	1999 Sep 20	3(1/1)	0.4	0.12	0.12	0.14		
4.....	U10	X095550.0+694046	...	39.51	361	1999 Sep 20	3(1/1)	0.4	0.11	0.11	0.13		
5.....	U10	X095550.0+694046	...	40.11	379	2000 Mar 11	3(1/1)	4.2	0.46	0.09	0.11		
6.....	U36	X132938.6+582505	NGC 5204	39.85	2028	2001 Jan 9	7(1/8)	0.6	0.43	0.10	0.12		
7.....	...	X140319.6–412258	NGC 5408	39.82	2885	2002 May 7	7(1/4)	0.7	0.32	0.08	0.11	NGC 5408 X-1	4, 5, 6
8.....	U43	X141312.2–652014	Circinus	39.30	356	2000 Mar 14	7(1/1)	0.4	0.12	0.12	0.15		
Comparison													
9.....	...	X095533.0+690033	M81	38.39	735	2000 May 7	7(1/1)	1.0	0.18	0.59	0.21	MF97	1, 4, 5, 6, 7

NOTES.—Col. (1): Observation number. Col. (2): Sample source name if the same as the one in Table 2. Col. (3): X-ray positions (J2000.0). Col. (4): Host galaxy. Col. (5): Approximate observed luminosity in units of erg s^{-1} in the energy band 0.3–8.0 keV, derived from the count rate of the reprocessed data; we assumed a PL model with $\Gamma = 1.8$ and a Galactic absorption column. Col. (6): Observation ID. Col. (7): Date of observation start. Col. (8): CCD number where the object is located and subarray values (in parentheses). Col. (9): Off-axis angle of the source in arcminutes. Col. (10): Count rate in s^{-1} . Col. (11): Reduced count rate calculated for pileup estimations explained in § 2.4; this takes into account the off-axis angle in col. (9) and the subarray values in col. (8). Col. (12): Pileup estimation based on the reduced count rate in col. (11). Col. (13): Common names from the literature in col. (14) (see Table 2 for common names and references for objects listed in col. [2]). Col. (14): References.

REFERENCES.—(1) Liu & Bregman 2005; (2) Colbert & Ptak 2002; (3) Goad et al. 2006; (4) Feng & Kaaret 2005; (5) Liu & Mirabel 2005; (6) Swartz et al. 2004; (7) Swartz et al. 2003.

the other hand, a well-known and previously studied ULX, IXO 85 (C22), is excluded from our ULX sample because the *Chandra* luminosity falls just below our ULX limit.

Examining the galaxy properties in Table 2, we find that most of our sample objects reside in spiral or irregular (merger) galaxies and are preferentially located in spiral arms and star-forming regions. Our galaxy sample includes two mergers (NGC 520 and the Antennae) and four early-type galaxies (NGC 2681, NGC 4125, M87, and Cen A). We see little difference between the ULX locations in their host galaxies in general and the locations of the comparison sources. The two groups also tend to have similar de-projected offsets from the centers of their galaxies. Two ULXs (U2 and U14) are associated with the nucleus of their host galaxies (M33 and NGC 3310, respectively), but with no evidence of AGN activity. We do not have enough detailed information on these sources to know what fraction are known low-mass X-ray binaries (LMXBs) and high-mass X-ray binaries (HMXBs). The identification of the optical counterparts would require sensitive optical imaging. Based on their location in the host galaxy, we can only say that most of our sources in both samples are consistent with being HMXBs.

2.6. Background Contamination

Here we estimate potential sample contamination from additional background objects that have not already been clearly identified. We use the $\log N$ - $\log S$ function from X-ray deep-field surveys to estimate the fraction of additional background objects based on our source fluxes and galaxy distances. For our sample criteria, we construct two flux limits, FL and FC. FL is the flux of a source with a specific luminosity: $10^{39.0} \text{ erg s}^{-1}$ for ULXs and $10^{38.3} \text{ erg s}^{-1}$ for our comparison sample. FC is the flux of a source that provides 1000 counts in its spectrum for the longest exposure time obtained for each galaxy. Assuming a PL model with $\Gamma = 1.8$ and using the Galactic value of absorption corresponding to the location of the center of each galaxy on the sky, we calculate FL and FC for all of the 286 galaxies in our original list (see § 2.1). The final flux limit for each galaxy to compare

with $\log N$ - $\log S$ is the largest of the two fluxes, FC or FL. All of our measured fluxes are above $10^{-14} \text{ erg cm}^2 \text{ s}^{-1}$, which corresponds to an ACIS count rate of $\sim 10^{-3} \text{ s}^{-1}$.

To make a background estimate we also need to account for the size of the detectors on the sky compared to the projected sizes of the galaxies. The area of each galaxy in deg^2 is first calculated within the D_{25} ellipse. Most observations are done in ACIS imaging mode, with detector areas of $\sim 0.117 \text{ deg}^2$ for both ACIS-I and ACIS-S. Data can be extracted from specific CCD chips, and some observations are only in subarray mode with a significantly smaller exposed area. The disparity between the sizes of the galaxies and the detector coverage can affect our background estimates. For the nearest galaxies, their size on the sky is larger than or comparable to the size of the ACIS detectors. Naturally, if the projected area of the galaxy is larger than or comparable to the size of the detector, these galaxies will provide the largest estimated contributions to the background counts. We therefore account for the fractional coverage of the 13 largest galaxies by overplotting the CCDs and estimating the coverage fraction. These 13 galaxies (out of the original 286) contribute 65% to the total estimate of the contamination. For the remaining galaxies we use the D_{25} ellipse area.

We use the $\log N$ - $\log S$ function from two separate surveys to obtain flux estimates. The popular *ROSAT* deep survey in the Lockman Hole (Hasinger et al. 1998) gives $\log N$ - $\log S$ for the flux interval 10^{-15} to $10^{-13} \text{ erg cm}^2 \text{ s}^{-1}$, in the range 0.5–2 keV. We apply a scale factor of 0.38 for our 0.3–8 keV band, obtained using the absorbed PL model with $\Gamma = 1.8$. The *Chandra* Multi-wavelength Project (ChaMP) serendipitous survey (Kim et al. 2004) contains a larger sample and covers a wide area ($\sim 14 \text{ deg}^2$). It uses the same soft X-ray band as the *ROSAT* deep survey, but the slope of the $\log N$ - $\log S$ function is shallower at the high end.

For ULXs, the *ROSAT* and ChaMP surveys predict no more than three or five spurious sources, respectively. For our lower luminosity objects, the prediction is one or two spurious sources. The survey estimates are compatible given large errors due to poor sampling at the high-flux end. Thus, no more than approximately 1

TABLE 2
PROPERTIES OF SAMPLE OBJECTS

No. (1)	Name (2)	Position (3)	Galaxy (4)	Dist. (5)	$N_{\mathrm{H}}^{\mathrm{gal}}$ (6)	ObsID (7)	CCD (8)	$\log L_{\mathrm{X}}$ (9)	Date (10)	Exp. (11)	Counts (12)	θ_{off} (13)	Alternate Names (14)	Location (15)	Refs. (16)
ULX Sample															
1.....	U1	X012435.2+034731	NGC 520	29.6	3.3	2924	7(1/1)	40.1	2003 Jan 29	49.3	1037.2 ± 32.3	2.1	Source 11	DB	1
2.....	U2	X013350.9+303938	M33	0.7	5.69	787	7(1/4)	39.0	2000 Jan 11	9.3	26733.4 ± 163.6	8.9	M33 X-8	N	2, 3, 4, 5
3.....		X013350.9+303939			5.69	2023	7(1/1)	38.9	2001 Jul 6	88.8	171301.0 ± 415.2	12.5			
5.....	U3	X022231.4+422024	NGC 891	9.6	8.12	794	7(1/1)	39.4	2000 Nov 1	50.9	1977.4 ± 44.5	1.7	NGC 891 X-4	DB	6, 7
6.....	U4	X024238.9-000055	M77	15.2	3.54	344	7(1/1)	39.7	2000 Feb 21	47.4	1524.8 ± 39.6	0.7		A	8
7.....	U5	X031820.0-662911	NGC 1313	3.7	3.96	3550	2(1/1)	40.1	2002 Nov 9	14.6	10486.7 ± 102.7	6.0	IXO 7, XMM1, NGC 1313 X-1	B	9, 64, 10, 11
8.....	U6	X034555.7+680455	IC 342	3.9	29.39	2916	7(1/8)	39.5	2002 Apr 29	9.3	2033.6 ± 45.1	0.5	IXO 22, IC 342 X-7, XMM1, X-1	A	9, 12, 64, 6, 13, 26, 47
9.....		X034555.6+680456			29.39	2917	7(1/8)	39.5	2002 Aug 26	9.9	2191.8 ± 46.8	0.6			
10.....	U7	X073625.5+653540	NGC 2403	4.2	4.17	2014	7(1/1)	39.2	2001 Apr 17	35.6	5364.2 ± 73.3	2.7	Source 21, NGC 2403 X-1, XMM1	A	14, 6, 64
11.....	U8	X085333.7+511930	NGC 2681	13.3	2.48	2061	7(1/1)	39.2	2001 May 2	79.0	1105.9 ± 33.3	1.3	NGC 2681 PSX-3	D	15
12.....	U9	X095546.5+694040	M82	5.2	4.02	361	3(1/1)	39.0	1999 Sep 20	33.3	1174.2 ± 34.9	0.8	Source 9	SF	16
13.....	U10	X095550.1+694048			4.03	378	3(1/1)	40.0	1999 Dec 30	4.1	1404.7 ± 38.0	4.0	Source 7, M82 X-1	SF	16, 17, 29
14.....	U11	X095551.0+694045			4.03	2933	7(1/1)	39.2	2002 Jun 18	18.0	1595.2 ± 41.3	0.6	Source 5	SF	16
15.....	U12	X095551.1+694043			4.03	361	3(1/1)	39.1	1999 Sep 20	33.3	1353.6 ± 39.1	0.4	Source 4	SF	16
16.....	U13	X103843.3+533102	NGC 3310	18.7	1.12	2939	7(1/2)	39.7	2003 Jan 25	47.2	1003.9 ± 31.7	0.3	IXO 38, NGC 3310 ULX2, X-3	A	9, 18, 19
17.....	U14	X103845.9+533012			1.11	2939	7(1/2)	39.8	2003 Jan 25	47.2	1541.8 ± 41.0	0.6	NGC 3310 X-1, X1	N	6, 18
18.....	U15	X103846.0+533004			1.11	2939	7(1/2)	39.8	2003 Jan 25	47.2	1221.6 ± 35.7	0.7		SF	
19.....	U16	X111126.0+554017	M108	14.1	0.78	2025	7(1/1)	39.4	2001 Sep 8	59.4	1278.9 ± 35.9	2.8	Source 26	D	20
20.....	U17	X112015.8+133514	NGC 3628	7.7	2.22	2039	7(1/1)	39.3	2000 Dec 2	58.0	2995.8 ± 54.8	0.9	IXO 39	DB	9, 21
21.....	U18	X120151.4-185225	NGC 4038/4039	21.7	3.95	3040	7(1/1)	39.7	2001 Dec 29	69.0	1009.7 ± 31.9	0.9	Source 11	AM	22, 23
22.....		X120151.3-185225			3.95	3043	7(1/1)	39.8	2002 Apr 18	67.1	1377.4 ± 37.3	1.0			
23.....		X120151.3-185225			3.95	3041	7(1/1)	39.8	2002 Nov 22	72.9	1491.4 ± 38.8	0.9			
24.....	U19	X120152.1-185134			3.95	315	7(1/1)	39.9	1999 Dec 1	72.2	1984.1 ± 44.6	1.6	Source 16	AM	22, 23
25.....		X120152.1-185133			3.95	3040	7(1/1)	39.8	2001 Dec 29	69.0	1587.0 ± 40.0	0.9			
26.....		X120152.1-185133			3.95	3042	7(1/1)	39.8	2002 May 31	67.3	1474.9 ± 38.5	1.6			
27.....		X120152.1-185133			3.95	3041	7(1/1)	39.8	2002 Nov 22	72.9	1491.8 ± 38.8	0.8			
28.....	U20	X120155.6-185215			3.96	315	7(1/1)	39.8	1999 Dec 1	72.2	1344.0 ± 37.3	1.9	Source 42	AM	22, 23
29.....	U21	X120156.4-185158			3.96	315	7(1/1)	39.7	1999 Dec 1	72.2	1307.1 ± 36.2	1.6	Source 44	AM	22, 23
30.....		X120156.5-185157			3.96	3040	7(1/1)	39.7	2001 Dec 29	69.0	1264.8 ± 35.6	0.4			
31.....	U22	X120807.5+651028	NGC 4125	18.1	1.82	2071	7(1/1)	39.5	2001 Sep 9	64.2	1051.7 ± 33.0	0.4		E	68
32.....	U23	X123030.6+414142	NGC 4485	9.3	1.78	1579	7(1/1)	39.6	2000 Nov 3	19.5	1450.1 ± 38.1	2.6	IXO 62, NGC 4485 X-1	A	9, 6, 24
33.....	U24	X123049.2+122604	M87	17.1	2.54	2707	7(1/1)	39.3	2002 Jul 6	98.7	1064.6 ± 36.8	3.1		E	68
34.....	U25	X123551.7+275604	NGC 4559	9.7	0.82	2026	7(1/4)	39.9	2001 Jan 14	9.4	1434.4 ± 37.9	0.6	IXO 65, NGC 4559 X-1, X7	D	9, 6, 25, 26, 48
35.....		X123551.7+275604			0.82	2027	7(1/4)	40.1	2001 Jun 4	10.7	2093.2 ± 45.8	0.6			
36.....	U26	X123558.6+275742			0.8	2027	7(1/4)	39.8	2001 Jun 4	10.7	1300.9 ± 36.1	2.9	IXO 66, NGC 4559 X-4, X10	DB	9, 6, 25, 26, 48
37.....	U27	X123617.4+255856	NGC 4565	16.4	1.31	3950	7(1/1)	39.8	2003 Feb 8	59.2	2146.5 ± 46.5	2.0	IXO 67, NGC 4565 ULX4	B	9, 27, 28
38.....	U28	X123740.3+114728	NGC 4579	20.3	2.52	807	7(1/4)	40.1	2000 May 2	33.9	1654.6 ± 40.7	1.3	NGC 4579 X-1	D	30
39.....	U29	X124155.6+323217	NGC 4631	6.9	1.29	797	7(1/1)	39.2	2000 Apr 16	59.2	3223.1 ± 56.8	0.5	IXO 68, NGC 4631 X-1, XMM1	SF	9, 6, 64

TABLE 2—*Continued*

No. (1)	Name (2)	Position (3)	Galaxy (4)	Dist. (5)	$N_{\mathrm{H}}^{\mathrm{gal}}$ (6)	ObsID (7)	CCD (8)	$\log L_X$ (9)	Date (10)	Exp. (11)	Counts (12)	θ_{off} (13)	Alternate Names (14)	Location (15)	Refs. (16)
40.....	U30	X125053.3+410714	M94	4.3	1.44	808	7(1/4)	39.0	2000 May 13	47.4	4472.6 ± 70.8	0.7	NGC 4736 X-1	DB	30
41.....	U31	X130521.9–492827	NGC 4945	5.2	14.94	864	7(1/1)	39.1	2000 Jan 27	49.1	2983.5 ± 54.9	1.3	NGC 4945 XMM4	DB	64
42.....	U32	X130532.9–492734			14.84	864	7(1/1)	39.1	2000 Jan 27	49.1	2797.7 ± 53.2	0.7	NGC 4945 X-2, XMM1	DB	31, 64
43.....	U33	X131519.5+420302	NGC 5055	7.2	1.3	2197	7(1/1)	39.8	2001 Aug 27	28.0	2354.6 ± 48.6	6.0	IXO 74, NGC 5055 X-2	D	9, 6
44.....	U34	X132507.4–430410	Cen A	4.9	8.41	316	3(1/1)	39.0	1999 Dec 5	35.7	1108.8 ± 33.9	9.2	IXO 75	E	9, 32
45.....		X132507.5–430410			8.41	962	1(1/1)	39.3	2000 May 17	36.5	2556.1 ± 50.7	5.5			
46.....	U35	X132519.8–430317			8.4	316	3(1/1)	39.2	1999 Dec 5	35.7	2124.0 ± 46.4	7.1	IXO 76	E	9, 32, 63
47.....	U36	X132938.6+582506	NGC 5204	4.8	1.38	2029	7(1/8)	39.4	2001 May 2	9.0	1498.1 ± 38.7	0.6	IXO 77, NGC 5204 X-1, XMM1	SF	9, (6, 18, 26, 33, 34), 64
48.....	U37	X133719.8–295349	M83	4.7	3.69	793	6(1/1)	39.0	2000 Apr 29	51.0	2419.2 ± 49.2	2.7	IXO 82, H30, XMM1	D	9, 35, 64
49.....	U38	X140304.0+542735	M101	5.4	1.15	4731	6(1/1)	39.2	2004 Jan 19	56.2	3213.5 ± 56.8	4.4	MF37, ULX2, H19, XMM-1, XMM2	A	39, 18, 37, 40, 64
50.....	U39	X140314.3+541806			1.15	5309	7(1/1)	39.0	2004 Mar 14	70.8	3889.1 ± 62.6	5.2	H25, P51, XMM-2, XMM1	A	37, 38, 40, 64
51.....		X140314.3+541806			1.15	4732	7(1/1)	39.0	2004 Mar 19	69.8	3902.9 ± 62.8	5.2			
52.....	U40	X140332.4+542103			1.15	934	7(1/1)	39.2	2000 Mar 26	98.2	9024.5 ± 95.1	3.8	M101 X5, H32, P98, ULX-1	A	18, 37, 38, 49, 65, 66, 60
53.....	U41	X140414.3+542604			1.15	934	3(1/1)	39.1	2000 Mar 26	98.2	3549.6 ± 59.7	10.6	IXO 83, ULX3, H45, XMM-3	D	9, 18, 37, 40
54.....		X140414.1+542603			1.15	4731	2(1/1)	38.8	2004 Jan 19	56.2	1085.5 ± 33.1	8.7			
55.....		X140414.2+542603			1.15	5300	3(1/1)	39.2	2004 Mar 7	52.1	2306.5 ± 48.2	11.0			
56.....		X140414.2+542603			1.15	5309	3(1/1)	38.9	2004 Mar 14	70.8	1592.2 ± 40.2	11.5			
57.....		X140414.2+542603			1.15	4732	3(1/1)	38.8	2004 Mar 19	69.8	1147.7 ± 34.3	11.5			
58.....	U42	X141310.1–652045	Circinus	3.7	59.7	356	7(1/1)	39.1	2000 Mar 14	23.1	1715.3 ± 41.4	0.9	CG X-2, source F	D	41, 42
59.....	U43	X141312.2–652014			59.92	365	7(1/8)	39.7	2000 Mar 14	5.0	1634.5 ± 40.5	0.4	CG X-1, source J	D	41, 42, 43
60.....	U44	X145358.9+033217	NGC 5775	22.4	3.51	2940	7(1/1)	39.9	2002 Apr 5	58.2	1324.2 ± 36.4	1.1		D	
61.....	U45	X203500.7+601131	NGC 6946	5.5	20.23	1043	7(1/1)	39.6	2001 Sep 7	58.3	8451.9 ± 92.1	4.8	MF16, NGC 6946 X-11, 58, X8	D	44, 6, 45, 46
62.....		X203500.8+601131			20.23	4404	7(1/1)	39.5	2002 Nov 25	30.0	3750.0 ± 61.3	2.9			
63.....	U46	X225724.7–410344	NGC 7424	11.5	1.33	3496	7(1/1)	39.7	2002 Jun 11	23.9	1370.8 ± 37.0	2.2	ULX2	A	67
64.....	U47	X225728.9–410212			1.32	3496	7(1/1)	39.7	2002 Jun 11	23.9	1331.9 ± 36.5	0.5	ULX1	D	67
Comparison Sample															
1.....	C1	X001528.9–391319	NGC 55	1.3	1.74	2255	0(1/1)	38.43	2001 Sep 11	59.4	8553.3 ± 92.5	3.7	Source 7, 6, N55	D	50, 7, 51, 52
2.....	C2	X004238.5+411604	M31	0.7	6.66	1585	0(1/1)	38.30	2001 Nov 19	4.9	1806.5 ± 42.6	4.3	r2-26, source 35	DB	(53, 62), (54, 55, 56)
3.....		X004238.5+411604			6.66	2895	0(1/1)	38.39	2001 Dec 7	4.9	2130.6 ± 46.3	5.3			
4.....		X004238.6+411603			6.66	2896	1(1/1)	38.40	2002 Feb 6	4.9	2302.4 ± 48.1	5.2			
5.....		X004238.6+411604			6.66	2898	3(1/1)	38.48	2002 Jun 2	4.9	2384.3 ± 49.0	6.8			
6.....	C3	X004305.7+411703			6.74	1575	7(1/1)	38.27	2001 Oct 5	37.7	20560.8 ± 143.5	4.8		DB	62, 55
7.....	C4	X004722.6–252051	NGC 253	3.0	1.35	790	6(1/1)	38.36	1999 Dec 27	43.5	1022.2 ± 32.3	7.8	NGC 253 PSX-5, X21, XMM2	A	15, (57, 58, 61), 64
8.....	C5	X004733.0–251749			1.37	969	7(1/1)	38.67	1999 Dec 16	14.0	1150.2 ± 34.0	0.3	NGC 253 PSX-2, X33, XMM1		15, (57, 58), 64
9.....		X004733.0–251749			1.37	790	6(1/1)	38.85	1999 Dec 27	43.5	3246.7 ± 57.4	5.5			
10.....	C6	X004735.2–251512			1.39	790	6(1/1)	38.57	1999 Dec 27	43.5	1811.8 ± 42.6	3.8	NGC253 PSX-7, X36, XMM3	A	15, (57, 58), 64
11.....	C7	X073655.6+653541	NGC 2403	4.2	4.17	2014	7(1/1)	38.94	2001 Apr 17	35.6	2608.9 ± 51.1	1.1	Source 20, XMM3	DB	14, 64
12.....	C8	X073702.4+653935			4.18	2014	7(1/1)	38.72	2001 Apr 17	35.6	1600.1 ± 40.1	5.0	Source 1, NGC 2403 X-4, XMM4	A	14, 6, 64
13.....	C9	X122809.3+440508	NGC 4449	3.0	1.5	2031	7(1/1)	38.38	2001 Feb 4	26.6	1138.8 ± 33.8	2.3	NGC 4449 X-1, source 10	SF	6, 59
14.....	C10	X122817.8+440634			1.49	2031	7(1/1)	38.46	2001 Feb 4	26.6	1356.9 ± 36.9	1.7	NGC 4449 X-7, source 27	SF	6, 59

TABLE 2—*Continued*

No. (1)	Name (2)	Position (3)	Galaxy (4)	Dist. (5)	$N_{\mathrm{H}}^{\mathrm{gal}}$ (6)	ObsID (7)	CCD (8)	$\log L_X$ (9)	Date (10)	Exp. (11)	Counts (12)	θ_{off} (13)	Alternate Names (14)	Location (15)	Refs. (16)
15.....	C11	X124211.1+323236	NGC 4631	6.9	1.29	797	7(1/1)	38.73	2000 Apr 16	59.2	1104.0 ± 33.3	3.1	NGC 4631 PSX-1, XMM5	D	15, 64
16.....	C12	X125050.3+410712	M94	4.3	1.44	808	7(1/4)	38.51	2000 May 13	47.4	1349.6 ± 37.0	0.3	M94 X-4	DB	30
17.....	C13	X125052.7+410719			1.44	808	7(1/4)	38.79	2000 May 13	47.4	2553.3 ± 51.6	0.6	M94 X-3	DB	30
18.....	C14	X125053.1+410712			1.44	808	7(1/4)	38.83	2000 May 13	47.4	2782.9 ± 53.6	0.6	M94 X-2	DB	30
19.....	C15	X130518.5–492824	NGC 4945	5.2	14.96	864	7(1/1)	38.69	2000 Jan 27	49.1	1115.3 ± 33.8	1.8	NGC 4945 XMM3	DB	64
20.....	C16	X130538.1–492545			14.74	864	6(1/1)	38.94	2000 Jan 27	49.1	1535.9 ± 39.3	2.6	Source 3, NGC 4945 XMM2	D	31, 64
21.....	C17	X133659.5–294959	M83	4.7	3.69	793	7(1/1)	38.54	2000 Apr 29	51.0	1249.9 ± 35.5	3.9	H17, source 28, M83 XMM2	A	35, 36, 64
22.....	C18	X133700.9–295203			3.7	793	7(1/1)	38.58	2000 Apr 29	51.0	1341.9 ± 40.8	2.0	source 44	DB	36
23.....	C19	X133704.3–295404			3.72	793	7(1/1)	38.59	2000 Apr 29	51.0	1381.1 ± 37.2	1.0	H26, source 62	A	35, 36
24.....	C20	X133704.4–295122			3.69	793	7(1/1)	38.63	2000 Apr 29	51.0	1527.7 ± 39.2	2.2	H27, source 64, M83 XMM3	DB	35, 36, 64
25.....	C21	X140228.3+541627	M101	5.4	1.14	5322	6(1/1)	38.76	2004 May 3	64.7	1235.2 ± 35.3	5.6	M101 XMM4	D	64
26.....	C22	X203500.1+600908	NGC 6946	5.5	20.13	1043	7(1/1)	38.91	2001 Sep 7	58.3	1894.7 ± 43.6	3.4	IXO 85, NGC 6946 X-9, 56, X7	A	9, 6, 45, 46

NOTES.—Col. (1): Observation number. Col. (2): Source name. Col. (3): X-ray positions (J2000.0). Col. (4): Host galaxy. Col. (5): Galaxy distance from Tully (1988) in Mpc. Col. (6): Galactic absorption column in units of $10^{20} \mathrm{cm}^{-2}$. Col. (7): Observation ID. Col. (8): CCD number where the object is located and subarray values (in parentheses); the subarray value represents the fraction of the CCD actually used in the observation. Col. (9): Approximate observed luminosity in units of $\mathrm{erg s}^{-1}$ in the energy band 0.3–8.0 keV, derived from the count rate of the reprocessed data; we assumed a PL model with $\Gamma = 1.8$ and a Galactic absorption column. Col. (10): Date of observation start. Col. (11): Exposure time in ks. Col. (12): Net counts in the 0.3–8.0 keV energy band. Col. (13): Off-axis angle of the source in arcminutes; the values listed here and the subarray values in col. (8) were used when we rejected the piled-up sources. Col. (14): Common names from the literature in col. (16) (the names correspond to references in the same order, and references that use the same name are in parentheses; some papers do not give special names or the names are given using the coordinates, and these were not used). Col. (15): Location in the galaxy; abbreviations are A, spiral arm; D, disk; DB, disk or bulge; E, elliptical galaxy, no special location; SF, star-forming region; AM, arm in merger; N, nucleus. Col. (16): References.

REFERENCES.—(1) Read 2005; (2) Dubus & Rutledge 2002; (3) Colbert & Mushotzky 1999; (4) Foschini et al. 2004; (5) La Parola et al. 2003; (6) Roberts & Warwick 2000; (7) Read et al. 1997; (8) Smith & Wilson 2003; (9) Colbert & Ptak 2002; (10) Colbert et al. 1995; (11) Miller et al. 2004; (12) Kong 2003; (13) Sugihro et al. 2001; (14) Schlegel & Pannuti 2003; (15) Humphrey et al. 2003; (16) Matsumoto et al. 2001; (17) Strohmayer & Mushotzky 2003; (18) Liu & Bregman 2005; (19) Jenkins et al. 2004a; (20) Wang et al. 2003; (21) Strickland et al. 2001; (22) Fabbiano et al. 2001; (23) Zezas et al. 2002; (24) Roberts et al. 2002; (25) Vogler et al. 1997; (26) Roberts et al. 2004; (27) Foschini et al. 2004; (28) Wu et al. 2002; (29) Mucciarelli et al. 2006; (30) Eracleous et al. 2002; (31) Guainazzi et al. 2000; (32) Kraft et al. 2001; (33) Roberts et al. 2001; (34) Roberts et al. 2005; (35) Immler et al. 1999; (36) Soria & Wu 2003; (37) Wang et al. 1999; (38) Pence et al. 2001; (39) Matonick & Fesen 1997; (40) Jenkins et al. 2004b; (41) Bauer et al. 2001; (42) Smith & Wilson 2001; (43) Weisskopf et al. 2004; (44) Roberts & Colbert 2003; (45) Holt et al. 2003; (46) Lira et al. 2000; (47) Roberts et al. 2003; (48) Cropper et al. 2004; (49) Mukai et al. 2003; (50) Schlegel et al. 1997; (51) Roberts 1997; (52) Stobbart et al. 2004; (53) Kong et al. 2002; (54) Primi et al. 1993; (55) Kaaret 2002; (56) Barnard et al. 2003; (57) Vogler & Pietsch 1999; (58) Pietsch et al. 2001; (59) Summers et al. 2003; (60) Kuntz et al. 2005; (61) Tanaka et al. 2005; (62) Williams et al. 2004; (63) Ghosh et al. 2006; (64) Winter et al. 2006; (65) Mukai et al. 2005; (66) Kong & Di Stefano 2005; (67) Soria et al. 2006; (68) Swartz et al. 2004.

in 10 sources in our sample is likely a background object. In a practical sense this is an upper limit, as our estimate does not take into account the variable absorption column through each galaxy, which will attenuate the signal of any background sources shining through the galaxy (i.e., reduce their observed flux). This is especially important, as we have used surveys in the 0.5–2 keV band where absorption is strong. We also remind the reader that we have already identified and rejected two background quasars (§ 2.3).

3. SPECTRAL ANALYSIS

We grouped the spectra to have a minimum of 15 counts per energy bin for the energy range of 0.3–8.0 keV. All fits were performed using the Galactic absorbing column (as listed in Table 2), plus an intrinsic absorbing column for each galaxy. Galactic values were obtained with the COLDEN routine in CIAO, which provides a foreground N_H value at a given celestial position. We chose to define acceptable (or “good”) fits as those for which $\chi^2_\nu \leq 1.2$. Unless specified, all errors quoted are 90% confidence for one interesting parameter ($\Delta\chi^2 = 2.7$). For sources with multiple observations, the individual observations were first fitted separately, and then all observations were fitted together in XSPEC for the various purposes of our work. Simultaneous fits are used in the histograms and listed in the tables (e.g., Table 3), and individual fits are shown in some of the plots to demonstrate any variability in luminosity and spectral shape. For the simultaneous fits, the model parameters were constrained to the same value in XSPEC, and only the normalizations of model components were allowed to vary freely.

In a statistical sense, spectral fitting results can be strongly biased by the number of counts in each spectrum. To test for such biases between the ULX and the comparison samples, we constructed histograms of net counts in the spectra (see Fig. 1). For sources with multiple observations, we chose the observation that contained the largest number of counts in its spectrum (see Table 2) to represent in the histogram. The distributions of the number of counts for objects in the samples are similar. A Kolmogorov-Smirnov (K-S) test provides a probability of 0.86, indicating that we have no reason to reject the hypothesis that the distributions are identical in count space. Thus, the two samples have equal sensitivity to spectral features for our model fitting.

3.1. Single-Component Spectral Fits

Recent spectral analysis of *Chandra* ULX spectra shows that many are well fitted by simple models (e.g., Humphrey et al. 2003; Swartz et al. 2004). We therefore fit all the spectra with either an absorbed PL model or a multicolor disk blackbody (MCD) model, with absorption fixed at the Galactic value in XSPEC. To keep our results within physical bounds, we impose upper limits of $\Gamma \leq 10$ and $kT_{\text{in}} \leq 4$ keV, respectively. The results are listed in Table 3. Both the ULX and the comparison samples are generally well fitted by the absorbed PL model (66% and 50%, respectively, are good fits, as indicated in col. [5] of the table). For the absorbed MCD model, good fits comprise 45% and 50% of the samples, respectively.

The histograms in Figure 2 show the distributions of the photon index and inner disk temperature, normalized to allow for easy comparison. For the full sample, we find no significant difference between ULXs and lower luminosity objects. Luminosity dependences are presented in Figure 3. For objects that have multiple observations, all fit results are shown.

We have applied the K-S test and the T -test to the samples in different ways. The first row of Table 4 shows the results of the test applied to the total set of fits, while the second row is restricted

to the “good” fits, as defined in the first paragraph of § 3. All calculated probabilities are higher than a 5% significance level, confirming that there are no significant differences when comparing the distributions or their means. We note that the derived probabilities differ, in some cases significantly, between the K-S test and the T -test, which is an indication that the distributions plotted in Figure 2 are possibly derived from intrinsic samples that do not have normal distributions and/or that our sample sizes are small (such tests are usually more reliable when applied to large samples).

Even with these caveats, we find an interesting trend if we limit our sample further. When only considering the good fits, the disk temperatures are marginally higher for ULXs (at 1.8 keV, with a significance level of 7%–8%). If we further use Figure 3a to split the ULXs themselves into two groups, with a luminosity break at 5×10^{39} erg s $^{-1}$, then we find that the highest luminosity ULXs have significantly harder X-ray spectra than both the lower luminosity ULXs and the comparison sample (rows 3–6 of Table 4).

Our primary result from applying single-component models is that all of the highest luminosity ULXs that are well fitted by the PL model possess hard X-ray spectra ($\Gamma \leq 2$ and $kT_{\text{in}} \geq 1.3$ keV). The most luminous ULXs have harder spectra, and those that are less luminous have spectral shapes similar to the comparison sample. Not all of the high-luminosity ULXs have hard spectra, however, and so we have further defined a subsample of nine very luminous and hard ULXs (see Fig. 3a): U4, U5, U10, U11, U14, U18, U19 (with four observations), U20, and U43. These all have luminosities in excess of 5×10^{39} erg s $^{-1}$ and photon indices < 1.7 . This subclass is discussed further in the next sections.

3.2. Two-Component Spectral Fits

We next fit all spectra with the frequently used two-component model that consists of a MCD model plus a PL. Typical spectral states observed in black hole binaries and some ULXs (e.g., Kubota et al. 2001) include a soft (high) state, with a prominent blackbody component having $kT \sim 1$ keV plus a steep ($\Gamma \sim 2.5$) PL tail, or a hard (low) state, with the thermal component being generally cooler or nonexistent and most of the energy carried in a shallower PL ($\Gamma \sim 1.8$). We also mention the very high state (VHS), characterized by high luminosities, a steep PL ($\Gamma > 2.5$), a relatively cool disk, and sometimes X-ray quasi-periodic oscillations (QPOs; see McClintock & Remillard 2006 for a detailed description).

We note that these spectral states have been traditionally measured in the 2–20 keV energy band and therefore may not be recognized easily in the *Chandra* band. For example, in the high state the PL component would be completely absent in our 0.3–8 keV band. Also, one of the most important signatures expected from an IMBH is a cool accretion disk component. The inner disk temperature in the MCD model scales with the black hole mass as $\propto M^{-1/4}$. For typical values of $kT_{\text{in}} \sim 1$ keV for a black hole binary with $10 M_\odot$ in the high state, we would expect cool disks with $kT_{\text{in}} \sim 0.1$ –0.3 keV. A number of ULXs with high-quality spectra from *Chandra*, *XMM-Newton*, and *RXTE* were found in the past few years to show soft components well fit by an MCD model in this range (see Miller et al. 2004).

To compare with published results and restrict model parameters enough to be useful for our purposes, we select a two-component model with fixed parameters. We assume inner disk temperatures of 0.25 or 1 keV to represent either a cool disk or a “normal” disk temperature, respectively (models PLMCD0.25 and PLMCD1.0). For ULXs, good fits are derived for 70% and 72% of the sample

TABLE 3
SINGLE-COMPONENT SPECTRAL FITS

SOURCE (1)	PL MODEL ^a					MCD MODEL ^a				
	Γ^b (2)	N_H^c (10^{21} cm^{-2}) (3)	Norm PL ^d (4)	Good Fits ^e (5)	χ^2/dof^f (keV) (6)	kT_{in}^g (10^{21} cm^{-2}) (7)	N_H^c (8)	Norm MCD ^h (9)	Good Fits ^e (10)	χ^2/dof^f (11)
ULX Sample										
U1.....	$2.56^{+0.23}_{-0.21}$	$4.8^{+0.8}_{-0.7}$	$7.8^{+12.0}_{-1.4} \times 10^{-5}$		68.6/53	$0.91^{+0.11}_{-0.10}$	$2.0^{+0.5}_{-0.4}$	$1.1^{+0.8}_{-0.4} \times 10^{-2}$		95.5/53
U2.....	2.11	2.4	5.5		2296.5/728	1.07	0.7	6.4		2792.4/728
U3.....	$1.94^{+0.14}_{-0.13}$	$7.6^{+0.9}_{-0.8}$	$1.4^{+1.0}_{-0.2} \times 10^{-4}$		70.0/100	$1.43^{+0.14}_{-0.12}$	4.8 ± 0.5	$5.9^{+2.4}_{-1.7} \times 10^{-3}$	G	91.8/100
U4.....	$0.81^{+0.14}_{-0.13}$	$5.3^{+1.2}_{-1.0}$	$3.8^{+10.3}_{-0.6} \times 10^{-5}$	G	93.1/86	(>3.61)	$5.6^{+0.8}_{-0.7}$	$1.9^{+2.7}_{-0.1} \times 10^{-4}$	G	94.7/86
U5.....	1.70 ± 0.05	4.5 ± 0.3	$2.2^{+0.2}_{-0.1} \times 10^{-3}$	G	304.7/289	$1.66^{+0.07}_{-0.06}$	2.5 ± 0.2	$6.9^{+1.1}_{-0.9} \times 10^{-2}$		353.6/289
U6.....	1.71 ± 0.09	3.7 ± 0.5	$6.0^{+0.7}_{-0.6} \times 10^{-4}$	G	215.7/215	$1.71^{+0.05}_{-0.12}$	1.3 ± 0.3	$1.6^{+0.5}_{-0.4} \times 10^{-2}$		259.8/215
U7.....	2.17 ± 0.07	4.2 ± 0.3	$4.4^{+0.8}_{-0.3} \times 10^{-4}$		260.3/182	$1.13^{+0.05}_{-0.04}$	2.1 ± 0.2	$3.9^{+0.7}_{-0.6} \times 10^{-2}$	G	172.6/182
U8.....	$1.88^{+0.16}_{-0.15}$	1.7 ± 0.4	$2.3^{+8.1}_{-0.3} \times 10^{-5}$	G	60.8/58	$1.21^{+0.13}_{-0.12}$	0.3 ± 0.2	$2.1^{+0.9}_{-0.7} \times 10^{-3}$	G	53.4/58
U9.....	$2.45^{+0.22}_{-0.20}$	$7.0^{+1.3}_{-1.2}$	$2.0^{+1.6}_{-0.4} \times 10^{-4}$	G	63.7/61	$1.11^{+0.12}_{-0.11}$	3.0 ± 0.8	$1.3^{+0.7}_{-0.5} \times 10^{-2}$		76.7/61
U10.....	$1.00^{+0.21}_{-0.20}$	$10.5^{+2.7}_{-2.4}$	$7.3^{+3.8}_{-1.9} \times 10^{-4}$	G	72.0/81	>3.42	$9.3^{+1.2}_{-1.1}$	<6.8	G	77.4/81
U11.....	$1.16^{+0.25}_{-0.23}$	$30.3^{+5.1}_{-4.4}$	$3.2^{+3.0}_{-1.0} \times 10^{-4}$	G	82.8/96	>3.23	$26.6^{+2.3}_{-2.1}$	$9.6^{+0.7}_{-0.5} \times 10^{-4}$	G	87.9/96
U12.....	$2.81^{+0.71}_{-0.63}$	$225.7^{+45.3}_{-38.8}$	$7.7^{+21.1}_{-5.3} \times 10^{-3}$	G	99.8/84	$1.89^{+0.58}_{-0.37}$	$179.2^{+30.5}_{-26.6}$	$2.1^{+4.2}_{-1.5} \times 10^{-2}$		101.8/84
U13.....	$2.55^{+0.27}_{-0.23}$	$2.3^{+0.5}_{-0.4}$	$4.4^{+10.8}_{-0.6} \times 10^{-5}$		69.9/50	$0.81^{+0.12}_{-0.11}$	$0.3^{+0.3}_{-0.2}$	$1.1^{+0.9}_{-0.5} \times 10^{-2}$		100.8/50
U14.....	$1.46^{+0.16}_{-0.14}$	$5.1^{+1.0}_{-0.9}$	$6.6^{+10.2}_{-1.0} \times 10^{-5}$	G	80.9/88	$2.07^{+0.34}_{-0.26}$	$3.3^{+0.6}_{-0.5}$	$1.2^{+0.8}_{-0.5} \times 10^{-3}$	G	87.2/88
U15.....	$1.78^{+0.18}_{-0.16}$	$6.5^{+1.3}_{-1.1}$	$7.3^{+12.9}_{-1.4} \times 10^{-5}$	G	76.8/70	$1.53^{+0.19}_{-0.16}$	$4.2^{+0.8}_{-0.7}$	$2.8^{+1.5}_{-1.0} \times 10^{-3}$	G	66.7/70
U16.....	$1.90^{+0.16}_{-0.15}$	$3.9^{+0.6}_{-0.5}$	$5.0^{+8.7}_{-0.7} \times 10^{-5}$	G	73.1/65	$1.30^{+0.14}_{-0.12}$	$2.1^{+0.4}_{-0.3}$	$3.2^{+1.4}_{-1.0} \times 10^{-3}$	G	65.9/65
U17.....	1.71 ± 0.11	7.6 ± 0.7	$1.5^{+0.9}_{-0.2} \times 10^{-4}$	G	149.7/149	$1.69^{+0.15}_{-0.14}$	$5.2^{+0.5}_{-0.4}$	$4.4^{+1.6}_{-1.2} \times 10^{-3}$		182.6/149
U18.....	1.79 ± 0.08	$4.0^{+0.3}_{-0.4}$	$4.7 \pm 0.4 \times 10^{-5}$		270.2/208	$1.48^{+0.10}_{-0.09}$	2.1 ± 0.2	$2.1^{+0.5}_{-0.4} \times 10^{-3}$		269.4/208
U19.....	$1.15^{+0.05}_{-0.02}$	0.6 ± 0.1	$2.2^{+0.2}_{-0.1} \times 10^{-5}$	G	391.6/345	$2.65^{+0.21}_{-0.22}$	<0.1	$2.6^{+0.8}_{-0.5} \times 10^{-4}$		416.7/345
U20.....	$1.22^{+0.13}_{-0.11}$	$0.2(<0.5)$	$1.8^{+6.8}_{-0.5} \times 10^{-5}$	G	62.2/75	$1.89^{+0.25}_{-0.24}$	<0.1	$6.1^{+4.0}_{-2.0} \times 10^{-4}$	G	84.6/75
U21.....	$1.97^{+0.11}_{-0.10}$	$1.4^{+0.3}_{-0.1}$	$(2.8 \pm 0.2) \times 10^{-5}$		203.7/133	1.16 ± 0.08	<0.2	$2.8^{+0.8}_{-0.6} \times 10^{-3}$		211.1/133
U22.....	$2.00^{+0.19}_{-0.17}$	0.7 ± 0.3	$2.1^{+7.9}_{-0.3} \times 10^{-5}$	G	48.5/53	$0.79^{+0.09}_{-0.08}$	<0.1	$9.0^{+4.3}_{-2.9} \times 10^{-3}$		84.3/53
U23.....	$1.80^{+0.14}_{-0.13}$	3.4 ± 0.5	$1.5^{+1.0}_{-0.2} \times 10^{-4}$		97.8/75	$1.41^{+0.14}_{-0.12}$	$1.7^{+0.4}_{-0.3}$	$8.3^{+3.3}_{-2.4} \times 10^{-3}$	G	80.5/75
U24.....	$2.52^{+0.28}_{-0.25}$	$1.9^{+0.5}_{-0.4}$	$2.1^{+9.6}_{-0.3} \times 10^{-5}$	G	59.1/66	$0.73^{+0.10}_{-0.09}$	0.3 ± 0.2	$8.3^{+5.8}_{-3.2} \times 10^{-3}$	G	70.4/66
U25.....	$2.25^{+0.11}_{-0.10}$	1.6 ± 0.2	$(2.5 \pm 0.2) \times 10^{-4}$		252.3/169	0.91	0.1	5.0		361.5/169
U26.....	$1.89^{+0.14}_{-0.13}$	2.5 ± 0.4	$2.2^{+1.1}_{-0.3} \times 10^{-4}$	G	70.1/67	$1.29^{+0.13}_{-0.11}$	$1.0^{+0.3}_{-0.2}$	$1.6^{+0.6}_{-0.4} \times 10^{-2}$	G	56.8/67
U27.....	2.00 ± 0.12	3.5 ± 0.4	$8.3^{+7.0}_{-0.9} \times 10^{-5}$	G	119.4/108	$1.20^{+0.10}_{-0.09}$	1.7 ± 0.2	$6.8^{+2.3}_{-1.7} \times 10^{-3}$		134.6/108
U28.....	1.88 ± 0.12	1.7 ± 0.3	$7.8^{+7.3}_{-0.8} \times 10^{-5}$	G	82.8/83	$1.33^{+0.12}_{-0.10}$	0.3 ± 0.2	$5.1^{+1.9}_{-1.4} \times 10^{-3}$	G	79.4/83
U29.....	1.90 ± 0.09	$2.7^{+0.3}_{-0.2}$	$1.0^{+0.6}_{-0.1} \times 10^{-4}$	G	158.7/135	$1.28^{+0.09}_{-0.08}$	$1.1^{+0.2}_{-0.1}$	$7.3^{+2.0}_{-1.6} \times 10^{-3}$		204.2/135
U30.....	$1.26^{+0.07}_{-0.06}$	0.2 ± 0.1	$8.6^{+4.2}_{-0.5} \times 10^{-5}$	G	166.7/184	$1.86^{+0.17}_{-0.12}$	<0.1	$3.0^{+0.8}_{-0.7} \times 10^{-3}$		265.5/184
U31.....	$2.27^{+0.11}_{-0.10}$	$5.7^{+0.6}_{-0.5}$	$2.5^{+1.0}_{-0.3} \times 10^{-4}$		188.8/142	1.13 ± 0.06	$2.9^{+0.4}_{-0.3}$	$1.9^{+0.5}_{-0.4} \times 10^{-2}$	G	154.9/142
U32.....	1.79 ± 0.11	$6.4^{+0.8}_{-0.7}$	$1.8^{+1.0}_{-0.2} \times 10^{-4}$	G	158.9/142	$1.60^{+0.13}_{-0.12}$	$3.9^{+0.5}_{-0.4}$	$5.6^{+1.8}_{-1.4} \times 10^{-3}$	G	153.1/142
U33.....	$2.50^{+0.14}_{-0.13}$	2.6 ± 0.3	$4.7^{+1.1}_{-0.5} \times 10^{-4}$		131.5/104	0.80 ± 0.06	0.7 ± 0.2	$1.3^{+0.5}_{-0.3} \times 10^{-1}$		179.1/104
U34.....	3.97	2.9	1.7		514.9/148	0.26	1.5	6.3		611.5/148
U35.....	$2.42^{+0.14}_{-0.13}$	2.6 ± 0.5	$2.3^{+1.0}_{-0.3} \times 10^{-4}$	G	93.9/103	0.95 ± 0.06	$0.2(<0.5)$	$3.3^{+1.0}_{-0.8} \times 10^{-2}$	G	84.2/103
U36.....	$3.11^{+0.22}_{-0.20}$	2.0 ± 0.3	$3.3^{+1.2}_{-0.3} \times 10^{-4}$	G	63.8/70	0.49 ± 0.05	$0.3^{+0.2}_{-0.1}$	$5.6^{+3.1}_{-1.9} \times 10^{-1}$		105.2/70
U37.....	2.43 ± 0.12	$1.9^{+0.4}_{-0.3}$	$1.3^{+0.8}_{-0.1} \times 10^{-4}$		150.6/106	0.90 ± 0.05	<0.1	$2.4^{+0.5}_{-0.4} \times 10^{-2}$		154.7/106
U38.....	1.86 ± 0.09	3.5 ± 0.4	$1.5^{+0.7}_{-0.1} \times 10^{-4}$	G	178.8/151	$1.46^{+0.09}_{-0.08}$	$1.5^{+0.3}_{-0.2}$	$6.8^{+1.7}_{-1.3} \times 10^{-3}$	G	149.1/151
U39.....	$2.17^{+0.08}_{-0.07}$	1.9 ± 0.1	$(1.0 \pm 0.1) \times 10^{-4}$		430.4/282	0.91	0.5	2.1		800.3/282
U40.....	$6.51^{+0.27}_{-0.22}$	3.9 ± 0.3	$1.9^{+0.4}_{-0.1} \times 10^{-4}$		315.6/84	0.16 ± 0.01	1.3 ± 0.1	$1.2^{+0.4}_{-0.3} \times 10^2$		257.2/84
U41.....	$3.49^{+0.04}_{-0.09}$	$4.2^{+0.1}_{-0.2}$	$(1.2 \pm 0.1) \times 10^{-4}$	G	466.7/399	0.57 ± 0.02	0.9 ± 0.1	$7.4^{+1.4}_{-1.1} \times 10^{-2}$		576.8/399
U42.....	$1.48^{+0.17}_{-0.16}$	$4.7^{+1.6}_{-1.4}$	$2.1^{+1.6}_{-0.4} \times 10^{-4}$		133.0/93	$2.25^{+0.40}_{-0.31}$	$2.0^{+1.0}_{-0.9}$	$2.8^{+1.9}_{-1.2} \times 10^{-3}$		140.7/93
U43.....	$1.46^{+0.17}_{-0.16}$	$5.3^{+1.7}_{-1.5}$	$1.0^{+0.7}_{-0.2} \times 10^{-3}$	G	76.8/87	$2.36^{+0.42}_{-0.31}$	$2.4^{+1.1}_{-1.0}$	$1.2^{+0.5}_{-0.5} \times 10^{-2}$	G	76.2/87
U44.....	$1.90^{+0.26}_{-0.24}$	$29.6^{+4.4}_{-3.9}$	$2.0^{+2.4}_{-0.6} \times 10^{-4}$	G	84.9/75	$1.96^{+0.32}_{-0.26}$	$22.0^{+2.7}_{-2.4}$	$2.3^{+1.7}_{-1.0} \times 10^{-3}$	G	82.9/75
U45.....	2.60	1.3	3.9		729.9/321	0.61	0.0	3.3		1436.3/321
U46.....	$2.29^{+0.18}_{-0.16}$	2.1 ± 0.4	$1.1^{+1.0}_{-0.1} \times 10^{-4}$	G	67.4/66	0.92 ± 0.10	0.3 ± 0.2	$2.0^{+1.0}_{-0.7} \times 10^{-2}$		99.4/66
U47.....	1.80 ± 0.14	2.3 ± 0.4	$10.0^{+8.5}_{-1.2} \times 10^{-5}$	G	62.5/70	$1.29^{+0.14}_{-0.12}$	0.9 ± 0.2	$7.4^{+3.3}_{-2.3} \times 10^{-3}$	G	69.9/70
Comparison Sample										
C1.....	$3.69^{+0.10}_{-0.09}$	4.7 ± 0.2	$10.0^{+1.1}_{-0.6} \times 10^{-4}$		377.3/157	0.55 ± 0.02	1.3 ± 0.1	$6.4^{+1.3}_{-1.0} \times 10^{-1}$		561.5/157
C2.....	1.49 ± 0.05	$1.3^{+0.3}_{-0.2}$	$8.1^{+0.6}_{-0.5} \times 10^{-4}$		590.0/447	1.86 ± 0.08	<0.1	$(2.2 \pm 0.3) \times 10^{-2}$		560.6/447
C3.....	2.85 ± 0.04	2.5 ± 0.1	$1.5^{+0.1}_{-0.0} \times 10^{-3}$		530.1/229	0.72 ± 0.01	0.3	$5.5^{+0.5}_{-0.4} \times 10^{-1}$		442.3/229
C4.....	$2.63^{+0.28}_{-0.24}$	2.0 ± 0.7	$7.4^{+14.0}_{-1.3} \times 10^{-5}$		66.9/51	0.74 ± 0.07	<0.2	$2.6^{+1.2}_{-0.8} \times 10^{-2}$		88.7/51
C5.....	$1.97^{+0.09}_{-0.08}$	3.8 ± 0.3	$(2.3 \pm 0.2) \times 10^{-4}$		249.4/204	1.34 ± 0.08	1.7 ± 0.2	$1.2^{+0.3}_{-0.2} \times 10^{-2}$		322.1/204
C6.....	2.27 ± 0.14	$5.3^{+0.7}_{-0.6}$	$1.7^{+1.0}_{-0.2} \times 10^{-4}$	G	102.9/90	$1.09^{+0.08}_{-0.07}$	2.7 ± 0.4	$1.5^{+0.5}_{-0.4} \times 10^{-2}$	G	95.5/90
C7.....	1.79 ± 0.10	$2.7^{+0.4}_{-0.3}$	$1.4^{+0.7}_{-0.1} \times 10^{-4}$	G	139.4/126	$1.40^{+0.11}_{-0.10}$	1.1 ± 0.2	$7.9^{+2.4}_{-1.9} \times 10^{-3}$	G	141.3/126
C8.....	$1.43^{+0.12}_{-0.11}$	$1.5^{+0.4}_{-0.3}$	$6.1^{+7.3}_{-0.7} \times 10^{-5}$	G	80.1/82	$1.88^{+0.25}_{-0.20}$	0.5 ± 0.2	$1.7^{+0.8}_{-0.6} \times 10^{-3}$	G	79.9/82
C9.....	$2.71^{+0.24}_{-0.21}$	1.9 ± 0.4	$8.3^{+9.9}_{-1.1} \times 10^{-5}$		73.2/54	0.62 ± 0.06	0.3 ± 0.2	$6.0^{+3.0}_{-1.9} \times 10^{-2}$		68.2/54
C10.....	$1.97^{+0.17}_{-0.16}$	$6.8^{+0.9}_{-0.8}$	$1.7^{+1.2}_{-0.3} \times 10^{-4}$		81.8/68	$1.38^{+0.17}_{-0.15}$	$4.1^{+0.6}_{-0.5}$	$7.7^{+4.1}_{-2.7} \times 10^{-3}$		103.5/68
C11.....	$2.69^{+0.34}_{-0.31}$	$42.9^{+6.3}_{-5.5}$	$5.2^{+4.9}_{-1.9} \times 10^{-4}$	G	46.5/62	$1.37^{+0.25}_{-0.15}$	$29.5^{+3.7}_{-3.3}$	$8.6^{+6.4}_{-3.7} \times 10^{-3}$	G	44.3/62
C12.....	$1.54^{+0.14}_{-0.13}$	0.6 ± 0.3	$3.0^{+7.5}_{-0.3} \times 10^{-5}$		85.3/70	$1.40^{+0.15}_{-0.14}$	<0.1	$2.2^{+0.9}_{-0.6} \times 10^{-3}$		100.9/70

TABLE 3—*Continued*

SOURCE (1)	PL MODEL ^a					MCD MODEL ^a				
	Γ^b (2)	N_H^c (10^{21} cm^{-2}) (3)	Norm PL ^d (4)	Good Fits ^e (5)	χ^2/dof^f (keV) (6)	kT_{in}^g (10^{21} cm^{-2}) (7)	N_H^c (8)	Norm MCD ^h (9)	Good Fits ^e (10)	χ^2/dof^f (11)
C13.....	$1.81^{+0.10}_{-0.09}$	0.6 ± 0.2	$6.2^{+5.6}_{-0.5} \times 10^{-5}$	G	96.5/119	1.08 ± 0.07	<0.1	$1.0^{+0.3}_{-0.2} \times 10^{-2}$	G	140.9/119
C14.....	$2.28^{+0.13}_{-0.12}$	1.0 ± 0.2	$7.9^{+6.9}_{-0.7} \times 10^{-5}$		147.6/108	0.65 ± 0.04	<0.1	$6.1^{+1.5}_{-1.3} \times 10^{-2}$		233.6/108
C15.....	$1.46^{+0.15}_{-0.14}$	$1.6^{+0.6}_{-0.5}$	$3.7^{+9.6}_{-0.5} \times 10^{-5}$	G	54.2/62	$1.77^{+0.28}_{-0.21}$	$0.5^{+0.4}_{-0.3}$	$1.2^{+0.7}_{-0.5} \times 10^{-3}$	G	48.2/62
C16.....	$1.86^{+0.15}_{-0.14}$	4.7 ± 0.9	$1.0^{+1.1}_{-0.2} \times 10^{-4}$	G	84.5/78	$1.51^{+0.16}_{-0.14}$	2.1 ± 0.6	$3.8^{+1.7}_{-1.2} \times 10^{-3}$	G	85.6/78
C17.....	$1.38^{+0.15}_{-0.14}$	$0.7^{+0.4}_{-0.3}$	$2.7^{+7.3}_{-0.3} \times 10^{-5}$	G	74.2/66	$1.81^{+0.27}_{-0.21}$	<0.1	$8.9^{+5.4}_{-3.1} \times 10^{-4}$		83.5/66
C18.....	$2.60^{+0.22}_{-0.20}$	2.0 ± 0.4	$5.5^{+9.4}_{-0.7} \times 10^{-5}$	G	65.8/74	$0.74^{+0.08}_{-0.07}$	$0.1(<0.3)$	$2.0^{+1.1}_{-0.7} \times 10^{-2}$	G	74.2/74
C19.....	$2.60^{+0.17}_{-0.16}$	3.8 ± 0.5	$8.6^{+9.3}_{-1.2} \times 10^{-5}$	G	74.5/67	0.87 ± 0.07	1.3 ± 0.3	$1.5^{+0.6}_{-0.4} \times 10^{-2}$	G	65.4/67
C20.....	$2.35^{+0.17}_{-0.15}$	3.0 ± 0.4	$7.5^{+8.4}_{-0.9} \times 10^{-5}$	G	79.3/75	0.93 ± 0.08	$0.9^{+0.3}_{-0.2}$	$1.3^{+0.5}_{-0.4} \times 10^{-2}$	G	83.0/75
C21.....	$2.30^{+0.18}_{-0.17}$	$5.6^{+0.9}_{-0.8}$	$8.4^{+11.4}_{-1.4} \times 10^{-5}$		85.6/63	$1.07^{+0.10}_{-0.09}$	$2.8^{+0.6}_{-0.5}$	$8.1^{+3.4}_{-2.4} \times 10^{-3}$	G	68.8/63
C22.....	$5.55^{+0.54}_{-0.45}$	$5.3^{+0.9}_{-0.8}$	$3.2^{+2.2}_{-0.5} \times 10^{-4}$		136.9/73	$0.26^{+0.03}_{-0.02}$	1.5 ± 0.5	$6.2^{+6.4}_{-3.0}$		163.0/73

^a Model names.^b Photon index for the PL model.^c Intrinsic absorbing hydrogen column density in units of 10^{21} cm^{-2} .^d Normalization constant for the PL model in units of photons $\text{keV}^{-1} \text{ cm}^{-2} \text{ s}^{-1}$ at 1 keV.^e The “good” fits, marked with a G, have $\chi^2_\nu \leq 1.2$.^f The χ^2 value for the fit and number of degrees of freedom.^g Temperature of the accretion disk at the inner radius for the MCD model.^h Normalization constant for the MCD model, in units of $R_{\text{in}}(\text{km})^2 \cos \theta / D(10 \text{ kpc})^2$, where $R_{\text{in}}(\text{km})$ is the inner radius of the accretion disk in units of km, $\cos \theta$ is the cosine of the inclination of the accretion disk from the line of sight, and $D(10 \text{ kpc})$ is the distance to the source in units of 10 kpc.

for PLMCD0.25 and PLMCD1.0, respectively. For the comparison sample, good fits are derived for 59% and 55% for PLMCD0.25 and PLMCD1.0, respectively. The ULXs do possess a higher percentage of good fits, but the difference is not statistically significant given our sample sizes. Figure 4 shows the distribution of photon indices. The shaded areas correspond to the subsample of nine ULXs with high luminosities and hard X-ray spectra as defined in the previous section (see Fig. 3a). In total, there is no significant difference between the ULXs and the lower luminosity sources. A K-S test for the difference between the distributions gives probabilities of 0.21 and 0.15 for PLMCD0.25 and

PLMCD1.0, respectively. However, the nine high-luminosity, hard ULXs clearly stand out. We note the very steep PL component in some spectra for the PLMCD1.0 model. These results correspond to the “nonstandard model” fits of Stobbart et al. (2006b).

We tried our two-component model with all parameters free (PLMCD model), but many parameters are not constrained. Moreover, as seen in Figure 5, the MCD component is very weak or practically nonexistent in many cases. The nine high-luminosity, hard ULXs have the weakest disk components, practically negligible. We note a very steep PL component in some spectra here, again indicating a nonstandard model. In these spectra the nonthermal component is soft and strongly absorbed, as shown by the large values of the flux ratios. Here we only comment further on specific results for spectra that were not well fitted with the simple models from § 3.1. Table 5 presents the PLMCD model results, and in Figure 6 we plot absorbed luminosities versus the photon index and disk temperature. The two samples do not show significant differences. Both samples possess cool disks, and there is no apparent correlation of the disk temperature with luminosity. The presence of this soft-disk component also causes the PL slopes to generally become steeper, compared to our single PL fits (Fig. 3).

The use of applying the F -test for an added spectral component (Protassov et al. 2002) is controversial, so we performed simulations to check the validity of the F -tests. For each spectrum we performed 500 simulations under a null model, a PL in this case. We first used the command `tc1out simpars`, available in XSPEC version 12, to generate simulated parameters from the original fits. This method uses simulations from a multivariate normal distribution based on the covariance matrix estimated in the original fit. The simulated F -test results are listed in parentheses in column (7) of Table 5. Any differences between the simulations and the classical F -test are small and generally fall within the errors corresponding to the number of simulations ($\sim 5\%$). The method described in Protassov et al. (2002) uses a complete Bayesian Monte Carlo simulation to sample from the posterior distribution (developed by van Dyk et al. 2001). Our method approximates the posterior distribution with a multivariate normal distribution

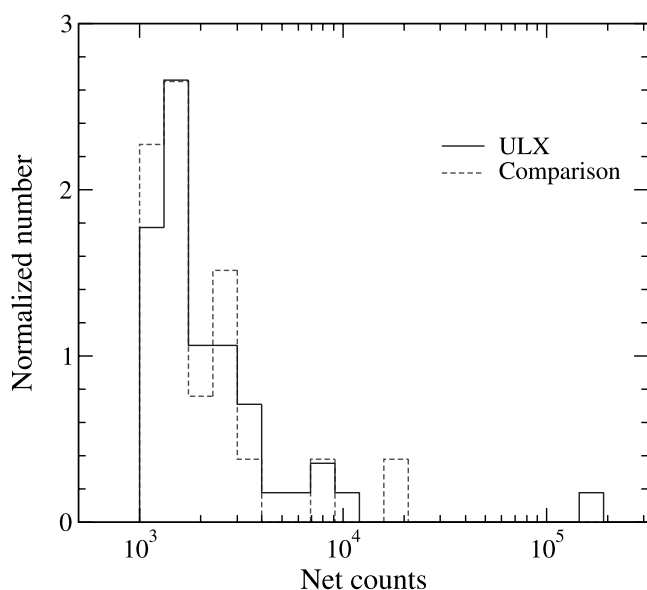


FIG. 1.—Normalized histograms of net counts for the ULX and comparison samples. For multiple observations we use the highest number of counts for each object. The histograms are normalized to unit area. The data with counts $>10^5$ are from one source: the long observation of M33 X-8 (U2). [See the electronic edition of the Journal for a color version of this figure.]

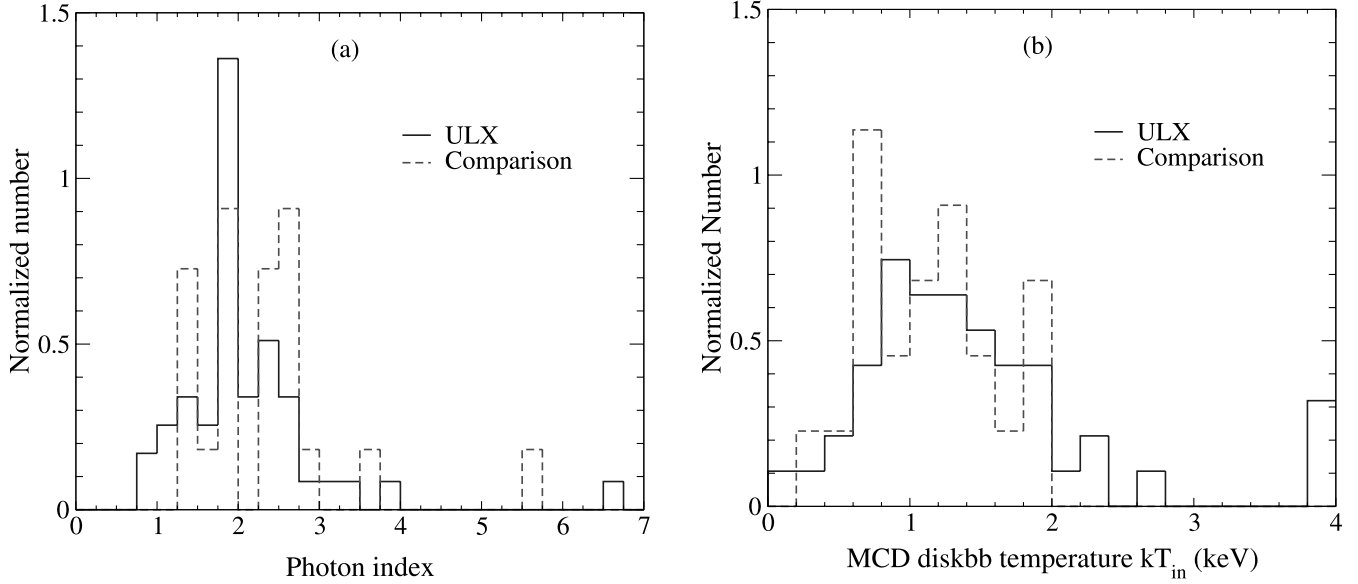


FIG. 2.—Normalized histograms from single-component fits. The histograms are normalized to have a unit area. (a) Photon index distribution from PL fits. (b) Inner disk temperature distribution from MCD fits. [See the electronic edition of the Journal for a color version of this figure.]

centered at the best-fit value. This is nevertheless better than just using the `fakeit` command on the original spectrum fitted with the null model (described by Protassov et al. [2002] as a “parametric bootstrap” and only valid when the parameters are very well constrained).

In conclusion, for the subsample of nine high-luminosity, hard ULXs (Fig. 3a and § 3.1) we recover the same result here; they tend to have significantly harder spectra. We also verify that they tend to possess small contributions from a thermal component. If such a component exists, it is practically undetectable with the *Chandra* data. We also find that cool disks (MCD with $kT_{\text{in}} \sim 0.1\text{--}0.3$ keV) are present with roughly equal probability for any given ULX and that cool-disk components extend below the standard ULX luminosity cutoff (10^{39} erg s $^{-1}$), down to our sample limit of $10^{38.3}$ erg s $^{-1}$.

4. SHORT-TERM VARIABILITY

Long-term flux variability from one observation to another, which is typically years, is very common in ULXs. Short-term flux variability, which we define here as that which can be detected within a single observation (hours), is less frequent and is not easily found with *Chandra*, probably due to limited sensitivity (i.e., not providing enough counts). Using the K-S test, Swartz et al. (2004) find that $\approx 15\%$ of our ULXs are variable at the 95% confidence level.

We extracted light curves for all sources, using three time bins: 3.24 (nominal frame time), 500, and 1000 s. To test for variability, we used the K-S statistic for the nominal frame time binning and the χ^2 test for the other two. Using the Monte Carlo method of Park et al. (2006) described in the previous section, we constructed

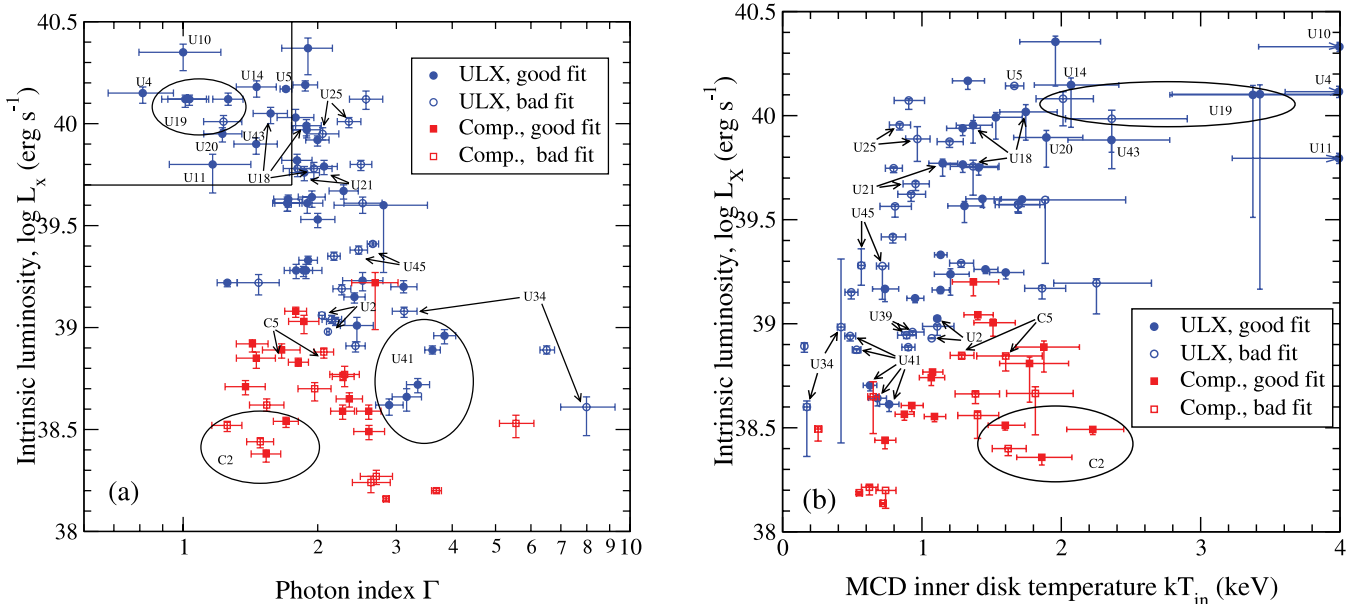


FIG. 3.—(a) Luminosity vs. photon index from PL model fits. (b) Luminosity vs. inner disk temperature using the MCD model. In the upper left corner of panel a we define a subsample of nine unique ULXs (U19 has four observations). They have luminosities $> 5 \times 10^{39}$ erg s $^{-1}$ and $\Gamma < 1.7$. For clarity, we label only the objects in this ULX subsample plus any objects with multiple observations.

TABLE 4
STATISTICAL TESTS FOR SINGLE-COMPONENT FITS

Samples	Γ (PL) ^a	K-S	<i>T</i> -Test	kT_{in} (MCD) ^b	K-S	<i>T</i> -Test	N_{H} (PL) ^c	K-S	<i>T</i> -Test	N_{H} (MCD) ^d	K-S	<i>T</i> -Test
ULX-All vs. Comp-All.....	2.11 ± 0.89 vs. 2.36 ± 0.92	0.155	0.324	1.44 ± 0.86 vs. 1.14 ± 0.47	0.457	0.126	21.51 ± 0.50 vs. 21.42 ± 0.42	0.74	0.49	21.18 ± 0.66 vs. 21.08 ± 0.55	0.84	0.59
ULX-GF vs. Comp-GF.....	1.88 ± 0.59 vs. 2.02 ± 0.50	0.731	0.479	1.81 ± 0.99 vs. 1.25 ± 0.37	0.075	0.078	21.56 ± 0.59 vs. 21.44 ± 0.50	0.70	0.55	21.33 ± 0.59 vs. 21.12 ± 0.63	0.32	0.39
ULX-HL vs. ULX-LL	1.70 ± 0.47 vs. 2.44 ± 1.05	0.011	0.005	1.89 ± 1.02 vs. 1.12 ± 0.50	0.015	0.001	21.54 ± 0.51 vs. 21.51 ± 0.50	0.75	0.86	21.25 ± 0.70 vs. 21.13 ± 0.64	0.47	0.58
ULX-HL-GF vs. ULX-LL-GF	1.52 ± 0.39 vs. 2.18 ± 0.57	0.020	0.001	2.26 ± 1.10 vs. 1.22 ± 0.27	0.022	0.012	21.58 ± 0.60 vs. 21.54 ± 0.61	0.52	0.87	21.52 ± 0.59 vs. 21.09 ± 0.53	0.43	0.11
ULX-HL vs. Comp-All	1.70 ± 0.47 vs. 2.36 ± 0.92	0.021	0.009	1.89 ± 1.02 vs. 1.14 ± 0.47	0.050	0.003	21.54 ± 0.51 vs. 21.42 ± 0.42	0.72	0.43	21.25 ± 0.70 vs. 21.08 ± 0.55	0.37	0.45
ULX-HL-GF vs. Comp-GF	1.52 ± 0.39 vs. 2.02 ± 0.50	0.111	0.009	2.26 ± 1.10 vs. 1.25 ± 0.37	0.023	0.008	21.58 ± 0.60 vs. 21.44 ± 0.50	0.48	0.53	21.52 ± 0.59 vs. 21.12 ± 0.63	0.18	0.15
ULX-LL vs. Comp-All	2.44 ± 1.05 vs. 2.36 ± 0.92	0.916	0.699	1.12 ± 0.50 vs. 1.14 ± 0.47	0.896	0.897	21.51 ± 0.50 vs. 21.42 ± 0.42	0.62	0.51	21.13 ± 0.64 vs. 21.08 ± 0.55	0.98	0.81
ULX-LL-GF vs. Comp-GF	2.18 ± 0.50 vs. 2.02 ± 0.50	0.727	0.460	1.22 ± 0.27 vs. 1.25 ± 0.37	0.637	0.843	21.54 ± 0.61 vs. 21.44 ± 0.50	0.93	0.64	21.09 ± 0.53 vs. 21.12 ± 0.63	0.97	0.90

NOTES.— Statistical comparison using the results from single-component fits. The samples compared in the first column are defined in § 3.1. The abbreviations are GF, good fits, with $\chi^2_{\nu} \leq 1.2$; HL, high-luminosity, ULXs with X-ray (absorbed) luminosity $L_X \geq 5.0 \times 10^{39} \text{ erg s}^{-1}$; and LL, low-luminosity, ULXs with $L_X \leq 5.0 \times 10^{39} \text{ erg s}^{-1}$. For each pair of samples in the first column we performed both a K-S and a *T*-test for the means, and we calculated the corresponding probabilities. The significant differences, with probabilities ≤ 0.05 , are shown in bold. High-luminosity ULXs have harder spectra than both low-luminosity ULXs and the comparison sample. There is also marginal evidence that ULXs show higher disk temperatures than the comparison sample if we only consider the good fits in both samples.

^a Average photon index in the PL model and 1 σ errors.

^b Average inner disk temperature in keV for the MCD model and 1 σ errors.

^c Average log H column density for the PL model and 1 σ errors in units of cm^{-2} .

^d The same for the MCD model.

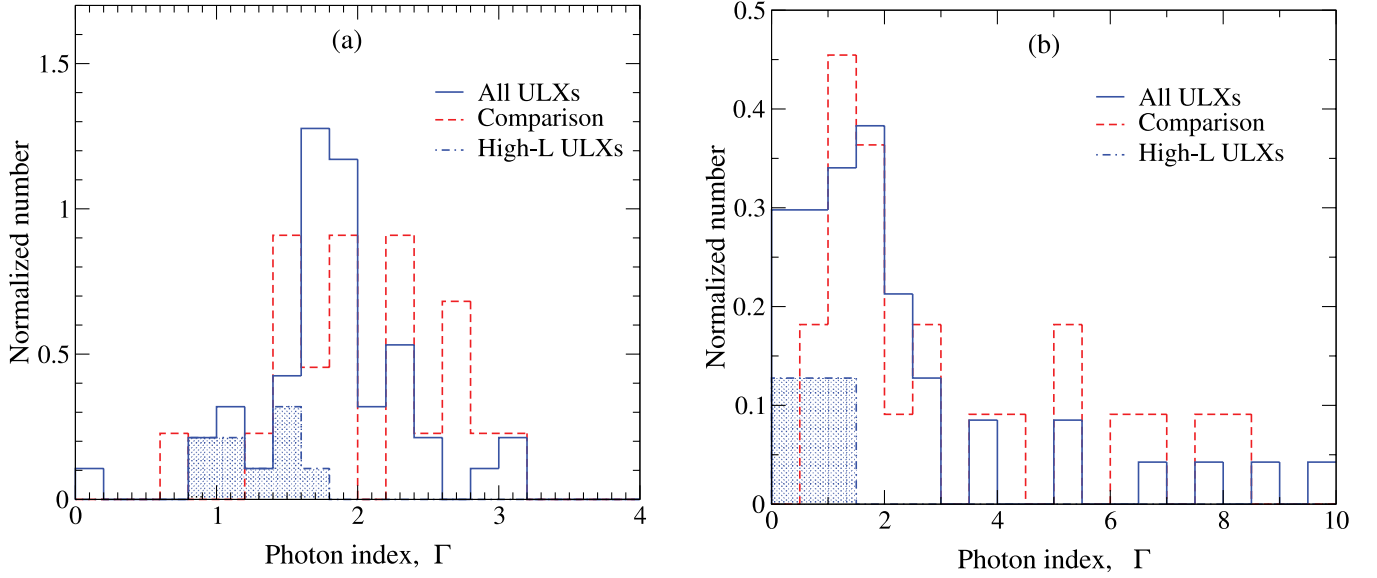


FIG. 4.—Histograms for photon indices from spectral fits with fixed inner disk temperatures, for ULX and lower luminosity samples, both normalized to unit area for easy comparison. We also show the high-luminosity, hard ULXs (filled blue regions). No significant difference is seen between ULXs and the comparison sample, but the high-luminosity ULXs are distinctly harder (i.e., flatter spectra). (a) Model PLMCD0.25. (b) Model PLMCD1.0

light curves for the hardness ratios for each variable source and looked for variations in hardness ratios and possible time lags between the three energy bands. We also constructed power spectra using the Leahy normalization.

We detect variability at 95% confidence in six ULXs for the longer time frames using the χ^2 test and no variability for the lower luminosity sample. Of these, three sources were previously known to be variable. These are M33 X-8, U34, and U40. Three other sources show variability. These are U14, U27, and C22, and the variability scale is similar to the exposure times of the obser-

vations (~ 40 ks). The K-S test identifies the same variable sources with the exception of U33 in NGC 5055, but it finds significant variability in two additional sources: U27 in NGC 4565 and a comparison source, C22 in NGC 6946. There are two periodicities of 707 s detected in U33 and U6 (NGC 1313 X-1) produced by the ACIS dither, which causes false periodic signals at 707 and 1000 s.

We conclude that six ULXs are intrinsically variable, which is consistent with the result obtained by Swartz et al. (2004) given the small size of the sample (47 sources). Only two sources (U2 and U40) show some energy variation, but no lag. Given the read-out time of the *Chandra* CCDs, variations on timescales shorter than ≈ 10 s cannot be detected, and features that could identify spectral states (QPOs or breaks in the power density spectra) are not readily detectable.

5. DISCUSSION

From our X-ray spectral comparison between ultraluminous X-ray sources (ULXs) and other X-ray point sources in nearby galaxies, we find an interesting subclass of nine ULXs that have unique properties compared to the other sources that are classified as ULXs. This subclass of ULXs also differs from the lower luminosity sample of X-ray point sources. While most of the ULXs we analyzed can plausibly be explained as scaled-up versions of Galactic black hole binaries, this particular subclass cannot. We discuss this subclass of ULXs followed by our general results, especially how our results relate to current evidence that supports the idea that ULXs host intermediate-mass black holes (IMBHs).

5.1. Luminous Hard (Flat-Spectrum) X-Ray ULXs

Our analysis has identified nine ULXs with very high luminosities and hard (flat) X-ray spectra. This sample (U4, U5, U10, U11, U14, U18, U19, U20, and U43) is shown in Figure 3a (upper left corner). These ULXs are all well fitted by a power-law (PL) model with a photon index of < 1.7 . More complex spectral fitting using two-component models reveals that when trying to add an accretion disk component (the MCD model described above), the contribution of this component to the total X-ray flux is very small, indicating that the relative contribution of emission

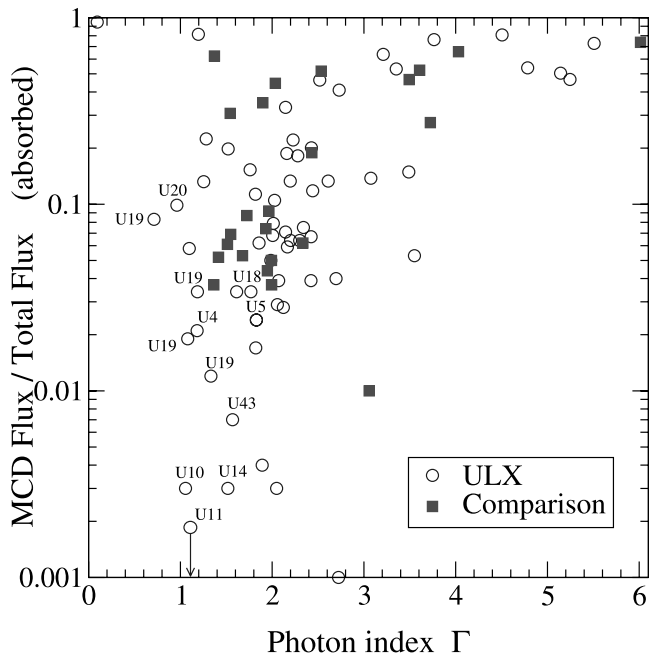


FIG. 5.—Ratio of MCD blackbody flux to the total flux (MCD fraction), plotted against photon index, using the free parameter model PLMCD. The fluxes are absorbed. For clarity, we only label the nine high-luminosity, hard ULXs as defined in Fig. 3a. These have both the hardest spectra and the lowest flux contribution from the MCD components. For U11 the fraction is below 0.001. [See the electronic edition of the *Journal* for a color version of this figure.]

TABLE 5
TWO-COMPONENT SPECTRAL FITS (MODEL PLMCD)

Source (1)	Γ^a (2)	kT_{in}^b (3)	Norm PL ^c (4)	Norm MCD ^d (5)	N_{H}^e (6)	$\Delta\chi^2/\text{Prob.}^f$ (7)	Good Fits ^g (8)	χ^2/dof^h (9)
ULX Sample								
U1.....	2.20 ^{+0.46} _{-0.69}	0.28 ^{+0.17} _{-0.12}	4.9 ^{+13.2} _{-3.2} $\times 10^{-5}$	1.1 ^{+1.0} _{-0.9}	5.0 ^{+2.6} _{-1.4}	3.1/0.69 (0.69)		65.5/51
U2.....	2.52 ^{+0.26} _{-0.15}	1.34 ^{+0.03} _{-0.06}	4.0 ^{+0.3} _{-0.1} $\times 10^{-3}$	1.4 $\times 10^{-1}$	2.5 ^{+0.2} _{-0.1}	397.9/>0.99 (>0.99)		950.9/424
U13.....	2.23 ^{+0.36} _{-0.04}	0.13 ^{+0.02} _{-0.02}	3.7 ^{+10.3} _{-1.3} $\times 10^{-5}$	<6.8 $\times 10^3$	5.0 ^{+2.0} _{-1.6}	27.4/>0.99 (>0.99)	G	42.5/48
U18.....	1.82 ^{+0.30} _{-0.17}	0.19 ^{+0.09} _{-0.05}	3.0 ^{+10.2} _{-0.5} $\times 10^{-5}$	<4.6 $\times 10^2$	6.7 ^{+4.1} _{-2.5}	8.0/0.94 (0.93)		68.1/51
U19.....	1.19 ^{+0.10} _{-0.14}	0.20 ^{+0.02} _{-0.19}	2.1 ^{+0.2} _{-0.3} $\times 10^{-5}$	1.1 ^{+0.9} _{-1.0}	<4.4	0.8/0.09 (0.18)		100.0/74
U25.....	1.76 ^{+0.28} _{-0.30}	0.19 ^{+0.08} _{-0.05}	1.6 ^{+1.0} _{-0.5} $\times 10^{-4}$	<5.8 $\times 10^2$	2.0 ^{+1.4} _{-0.8}	12.7/>0.99 (>0.99)		87.0/70
U25.....	2.16 ^{+0.19} _{-0.21}	0.13 ^{+0.03} _{-0.02}	3.0 ^{+1.0} _{-0.6} $\times 10^{-4}$	<9.8 $\times 10^3$	3.8 ^{+1.4} _{-1.0}	33.7/>0.99 (>0.99)		112.3/93
U33.....	2.34 ^{+0.23} _{-0.37}	0.20 ^{+0.13} _{-0.06}	4.0 ^{+1.4} _{-0.1} $\times 10^{-4}$	<4.7 $\times 10^2$	2.7 ^{+1.0} _{-0.7}	3.4/0.73 (0.71)		128.2/102
U34.....	3.21 ^{+0.75} _{-0.72}	0.10 ^{+0.02} _{-0.01}	9.6 ^{+22.2} _{-4.0} $\times 10^{-5}$	1.4 ^{+0.8} _{-1.2} $\times 10^5$	9.0 ^{+0.8} _{-1.1}	63.6/>0.99 (>0.99)	G	58.0/49
U34.....	2.73 ^{+0.19} _{-0.16}	0.11 \pm 0.01	2.2 ^{+0.7} _{-0.3} $\times 10^{-4}$	3.1 ^{+15.0} _{-2.4} $\times 10^4$	7.0 ^{+1.3} _{-1.1}	149.4/>0.99 (>0.99)		130.1/93
U37.....	3.35 ^{+0.43} _{-0.70}	1.14 ^{+0.15} _{-0.23}	1.0 ^{+2.7} _{-0.3} $\times 10^{-4}$	6.1 ^{+5.9} _{-2.4} $\times 10^{-3}$	2.4 ^{+2.7} _{-1.4}	11.2/0.98 (0.96)		139.4/104
U39.....	1.52 ^{+0.23} _{-0.29}	0.25 ^{+0.06} _{-0.05}	(4.8 \pm 1.6) $\times 10^{-5}$	3.7 ^{+8.1} _{-2.4}	2.1 ^{+0.5} _{-0.4}	54.9/>0.99 (>0.99)	G	142.8/139
U39.....	1.28 ^{+0.23} _{-0.25}	0.26 \pm 0.04	3.8 ^{+5.6} _{-1.1} $\times 10^{-5}$	3.6 ^{+5.5} _{-2.0}	2.1 \pm 0.4	83.9/>0.99 (>0.99)	G	148.5/137
U40.....	3.77 ^{+0.50} _{-0.42}	0.13 \pm 0.01	3.2 ^{+4.9} _{-0.9} $\times 10^{-5}$	4.6 ^{+3.2} _{-1.4} $\times 10^2$	1.8 ^{+0.3} _{-0.2}	169.5/>0.99 (>0.99)		146.2/82
U42.....	1.82	0.13	3.6	5.7	11.6	1.6/0.42 (0.47)		131.4/91
U45.....	2.43 ^{+0.12} _{-0.10}	0.12 \pm 0.01	3.5 ^{+0.6} _{-0.4} $\times 10^{-4}$	4.2 ^{+8.9} _{-2.4} $\times 10^3$	4.1 ^{+0.7} _{-0.5}	225.6/>0.99 (>0.99)	G	212.6/182
U45.....	2.28 ^{+0.18} _{-0.14}	0.13 \pm 0.02	2.9 ^{+0.9} _{-0.5} $\times 10^{-4}$	2.8 ^{+13.5} _{-2.0} $\times 10^3$	4.7 ^{+1.2} _{-0.9}	96.5/>0.99 (>0.99)		160.7/133
Comparison Sample								
C1.....	3.72 ^{+0.11} _{-0.08}	0.10 \pm 0.01	(1.2 \pm 0.1) $\times 10^{-3}$	1.3 ^{+3.3} _{-0.9} $\times 10^5$	8.9 \pm 0.4	137.6/>0.99 (>0.99)		239.7/155
C2.....	1.73 ^{+0.13} _{-0.18}	0.12 ^{+0.02} _{-0.01}	1.1 ^{+0.3} _{-0.1} $\times 10^{-3}$	1.3 ^{+12.6} _{-1.1} $\times 10^4$	4.7 ^{+1.6} _{-2.1}	8.5/0.95 (0.94)		144.4/108
C3.....	3.61 ^{+0.66} _{-0.38}	0.82 ^{+0.03} _{-0.04}	1.0 ^{+0.2} _{-0.1} $\times 10^{-3}$	2.1 ^{+0.5} _{-0.4} $\times 10^{-1}$	2.7 ^{+0.7} _{-0.4}	189.0/>0.99 (>0.99)		341.1/227
C4.....	2.43 ^{+0.42} _{-0.60}	0.18 ^{+0.17} _{-0.02}	6.1 ^{+14.2} _{-3.5} $\times 10^{-5}$	1.7 ^{+2.9} _{-1.6} $\times 10^1$	3.2 ^{+4.1} _{-2.0}	3.0/0.68 (0.65)		63.9/49
C5.....	1.96 \pm 0.19	0.20 ^{+0.07} _{-0.05}	2.3 ^{+1.3} _{-0.5} $\times 10^{-4}$	<7.0 $\times 10^2$	6.4 ^{+2.2} _{-1.6}	12.8/>0.99 (0.98)	G	164.7/139
C9.....	2.54(<3.49)	0.59 ^{+0.06} _{-0.14}	2.8 ^{+9.0} _{-2.8} $\times 10^{-5}$	5.2 ^{+11.6} _{-1.2} $\times 10^{-2}$	1.1 ^{+1.4} _{-0.9}	6.8/0.92 (0.90)		66.4/52
C10.....	1.99 ^{+0.17} _{-0.19}	0.18(<0.74)	1.8 ^{+0.4} _{-0.3} $\times 10^{-4}$	<8.5 $\times 10^1$	9.7 \pm 0.6	2.2/0.59 (0.60)		79.6/66
C12.....	1.42 ^{+0.23} _{-0.51}	0.31	2.5 ^{+7.5} _{-1.5} $\times 10^{-5}$	<1.4 $\times 10^{-1}$	0.6(<1.3)	1.2/0.39 (0.39)		84.0/68
C14.....	1.54 ^{+0.37} _{-0.52}	0.35 ^{+0.08} _{-0.07}	3.0 ^{+6.9} _{-1.6} $\times 10^{-5}$	4.3 ^{+5.6} _{-2.5} $\times 10^{-1}$	0.5 ^{+0.4} _{-0.3}	17.4/>0.99 (>0.99)		130.2/106
C22.....	3.49 ^{+0.62} _{-0.56}	0.12 \pm 0.02	8.9 ^{+36.8} _{-3.6} $\times 10^{-5}$	2.2 ^{+20.7} _{-1.7} $\times 10^3$	5.3 ^{+1.7} _{-1.1}	57.4/>0.99 (>0.99)	G	79.6/71

^a Photon index for the PL model.

^b Temperature of the accretion disk at the inner radius for the MCD model in keV.

^c Normalization constant for the PL model as in Table 3.

^d Normalization constant for the MCD model as in Table 3.

^e Intrinsic absorbing hydrogen column density in units of 10^{21} cm^{-2} .

^f F -test $\Delta\chi^2/\text{confidence levels}$ for the model PLMCD against the PL model alone. The values in parentheses are obtained from simulations; see § 3.2 for details.

^g The “good” fits are marked with a G as in Table 3.

^h The χ^2 value for the fit and number of degrees of freedom.

from the accretion disk to the total spectrum is small. This result is shown in Figure 5, where the ratio of the absorbed MCD component flux to the total flux is plotted against the PL index for the PLMCD model. By using absorbed fluxes we do not specifically show the absolute physical strength of the accretion disk component; however, the ratios measured in this way better indicate the significance of detecting a soft excess and are also less dependent on the modeling.

The nine members of the subclass of flat-spectrum ULXs are also very luminous, so they would seem to be the best candidates for hosting an IMBH based on a simple Eddington limit argument, which predicates that higher mass black holes are required to explain the most luminous accreting sources. The spectra of these ULXs resemble Galactic black holes in a hard state, but such spectral shapes are usually associated with a low-luminosity state in the case of Galactic black holes. If these ULXs are indeed accreting IMBHs in a low state (i.e., low/hard state), our result begs the question as to why we do not also see ULXs in a high state (high/soft state) with even higher luminosities. Moreover, if these are IMBHs in a low state, such a scenario implies very high mass

black holes ($>10^4 M_{\odot}$). The formation of such black holes is not easy to explain.

It seems more plausible that this subclass of hard (flat) and luminous ULXs is accreting sources in the PL-dominated very high state (VHS; McClintock & Remillard 2006), with an unusually weak soft X-ray component. A model that describes the properties of a hard PL with very little flux from the disk, at least in the *Chandra* band, is the “coupled disk-corona” model proposed by Done & Kubota (2006). In this model, the underlying accretion disk emission is distorted by a process that drains energy from the disk into the corona. In an extreme case, the inner disk emission could be almost completely Comptonized, and thus only the visible outer disk would contribute to the accretion disk component. As this obviously only appears at low temperatures it could be easily absorbed in some galaxies and also hard to detect. However, if our nine luminous ULXs are interpreted as stellar mass black hole systems in a high state, we would need to explain why their X-ray spectra are much harder compared to those observed in our Galaxy (McClintock & Remillard 2006), which have typical photon indices of 2.5 in the VHS. We would also need to

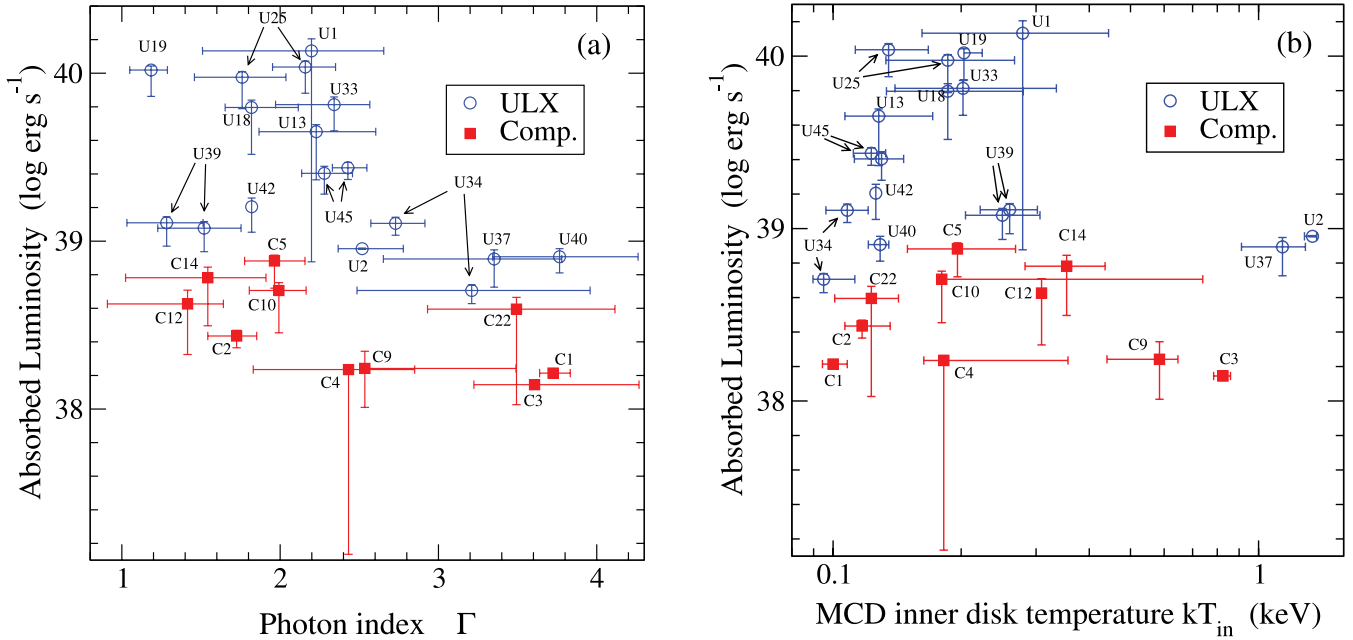


FIG. 6.— Absorbed luminosity scatter plots from the two-component spectral model with free parameters (PLMCD). We present results only for the spectra that did not provide acceptable fits with single-component models. (a) Photon index dependence. (b) Disk temperature dependence.

explain how such low-mass black hole systems could reach such high luminosities. A model that could explain both the flat non-thermal component and the weak soft-disk component in luminous ULXs has been proposed by Socrates & Davis (2006). Their ULX model shows that at super-Eddington accretion rates, in the inner region of the disk, magnetic fields in the corona can prevent strong winds; thus, the radiative efficiency is not reduced by photon trapping. The resulting spectrum is dominated by the coronal emission from the inner region, and the soft thermal component is generated only in the outer disk.

There is some direct evidence that high-luminosity states in ULXs correlate with the hardness of the PL tail in their spectra. Roberts et al. (2006) have shown this for a long *Chandra* observing campaign of NGC 5204 X-1 (our U36). These data were not available when we searched the archive. Roberts et al. (2006) found that the spectrum becomes harder as the flux increases over timescales of days to weeks. The model used was a Comptonized disk model, and the results showed a cool disk (~ 0.1 keV) and an optically thick corona. This model demonstrates that flux variations correlate with the corona temperature. Roberts et al. (2006) favor a stellar mass black hole interpretation for this ULX and suggest an unusual VHS, probably produced by extreme mass transfer from a massive star. We should note, however, that the PL slope is much steeper for NGC 5204 X-1 compared to the ULXs we discuss; therefore, it is much easier to interpret as a VHS in comparison to what is observed for Galactic black holes.

It is possible that our nine ULX spectra appear to be harder than they actually are due to the limited energy band covered by *Chandra*. If these objects intrinsically possess a break or curvature in their spectra, and the break occurs at an energy above the *Chandra* bandpass or where the sensitivity of *Chandra* falls off significantly, this might bias our modeling to measuring these sources as “hard.” This appears to be the case for one ULX in our sample, NGC 1313 X-1 (U5). Using *XMM-Newton* spectra, Stobbart et al. (2006a) found evidence for a break (or curvature) at 4.9 keV, with a photon index for the high-energy PL of 2.16 (much closer to the Galactic black hole). The authors show, how-

ever, that such breaks are easier to explain if ULXs contain stellar mass black holes rather than IMBHs. The curvature would be likely to originate in optically thick coronae. This theory would need to be tested for the remaining ULXs in our subclass by obtaining better quality spectra.

In conclusion, the subclass of nine ULXs with very high luminosities and hard (flat) X-ray spectra suggests a PL-dominated VHS, in line with recent models of stellar mass black hole systems in very high accretion states. The fact that the highest luminosity ULXs are explained more easily with such models argues strongly against IMBHs as the only explanation of ULXs.

5.2. Cool Disks and the IMBH Interpretation

Using our sample of ULXs and lower luminosity X-ray sources, we have found that the spectral signature of a cool accretion disk is not specific to ULXs. The results of the widely used PLMCD model (see Fig. 6 and Table 5) show that many sources in both samples have MCD components with low inner-disk temperatures. Cool disks have been used until recently as support for the IMBH interpretation. Our results show that this evidence is nowhere near as compelling when a large sample of ULXs is looked at properly.

In the standard accretion disk model, cool disks are not expected for stellar mass black holes accreting near their Eddington limits. The disk temperature scales with the black hole mass as $T_{\text{in}} \propto M^{-1/4}$ and is ~ 1 keV for stellar mass black holes. However, cool disks can be seen in a low (hard) state, because the temperature dependence on the accretion rate for standard disks is $T_{\text{in}} \propto \dot{M}^{-1/4}$ (McClintock & Remillard 2006; Miller et al. 2006). Cool disks have indeed been found in some non-ULX sources (e.g., Stobbart et al. 2006a). The authors note the similarity of these spectral fits with those typically used for ULXs. They also suggest that the soft excess in some cases could be otherwise explained by contamination from the host galaxy.

Most sources in Figure 6b possess low disk temperatures within both the ULX and comparison samples. Indeed, it is surprising that we do not see many states that are typical (high) states for

stellar mass black hole binaries with a prominent ~ 1 keV black-body component (although a hard PL tail would be difficult to discern in the limited *Chandra* bandpass). Only U2, U37, C3, and C9 show such high temperatures. We find a similar result from our cool-disk model (PLMCD0.25). A large fraction of our ULX spectra (70%) are well fitted by this model, but a significant number of lower luminosity objects (59%) are as well.

There are theoretical models that do not require the presence of an IMBH to explain a cool disk at high accretion rates and high luminosities (i.e., for ULXs). We already mentioned the model proposed by Done & Kubota (2006) that explains cool disks by a process of draining energy from the disk to launch an optically thick corona that obscures the hot inner regions of the disk. Freeland et al. (2006) developed microblazar models with magnetized jets that cause a transference of disk energy into the jet, thus making the disk fainter and cooling it at the same time. Other phenomenological models include the “dual thermal” model of Stobbart et al. (2006b), in which the soft excess comes from an optically thick outflow produced at high accretion rates (see King & Pounds 2003), which is seen in addition to a disk component with a temperature similar to those seen in stellar mass black hole binaries. This latter model was proposed to explain the alternate model of Stobbart et al. (2006b).

Done & Kubota (2006) found evidence of a Galactic black hole that supports the interpretation of ULXs as stellar mass black holes in a VHS. The microquasar XTE J1550–564 has a “strong VHS” (see also Kubota & Done 2004), where the disk temperature decreases with luminosity, reaching values of 0.3–0.4 keV. A similar behavior has been found for NGC 1313 X-2 (Feng & Kaaret 2007). This suggests a new type of VHS, a so-called ultra-luminous branch, which is very similar to the ULX spectra (Roberts 2007; Soria et al. 2007). In this interpretation, ULXs represent the high end of such an accretion state, with black hole masses up

to $100 M_{\odot}$ and accretion rates up to 20 times the Eddington limit. Forming black holes with such masses is much easier to explain than forming IMBHs. For example, Belczynski et al. (2004) showed that black holes with masses of $80 M_{\odot}$ or more can form through binary mergers. Soria & Wong (2006) suggested that black holes of up to $200 M_{\odot}$ could form by large-scale dynamical collapse of protoclusters in active regions in galaxies. These formation mechanisms are supported by the association between ULXs, star-forming regions, and colliding galaxies.

5.3. Conclusion

We have found that the highest luminosity ULXs tend to have the hardest X-ray spectra in the *Chandra* bandpass and are well fitted by a simple power-law model, without evidence for thermal accretion disk components. Such spectra are not consistent with current IMBH models but are more in line with current models of extreme very high states, or perhaps a new “ultraluminous state” (Roberts 2007), in stellar mass black holes.

Our work shows that cool accretion disks are not exclusive to the ULX class, suggesting that low-temperature IMBHs are not the only explanation for this phenomenon. In general, our results show that ULXs are likely to be composed of several distinct types of objects and that these types may extend into lower X-ray luminosity classes, such as classical Galactic black hole candidates and other objects in our comparison sample. Our conclusions provide another “nail in the coffin” for assumptions that ULXs are simply a class of accreting IMBHs.

No other specific properties have been found for the ULX group, except for spectral hardening at the highest luminosities. All these results suggest that ULXs are the highest luminosity end of stellar mass black hole binaries, with the largest black holes permitted by current formation mechanisms and/or accreting at super-Eddington rates.

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